Equatorial dynamics of hot Jupiters and warm Neptunes

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Overview

I. Observations and simulations

II. Warm Neptunes: breaking of the primitive equations

III. Hot Jupiters: spin-up of superrotation



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Hot Jupiters

First discovery: 51 Peg b, Mayor & Queloz 1995

Jupiter mass and radius

П D 2094 300					
Mass	Radius	Orbital period	Semi major axis		
$0.7 M_{J} = 222 M_{T}$	95000 km = 1.36 R _J	3.5 days	0.04 AU		

UD 2001506

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Short orbits: tidally locked

Day side – night side temperature difference: ~1000K

Warm Neptunes

First characterized: Corot 7b, Leger et al. 2009

GJ 1214b

Mass	Radius	Orbital period	Semi major axis
6.5 M _T	2.7 M _T	1.5 days	0.015 AU

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Warm Neptunes or super Earths?

First atmospheric studies: GJ 1214b. Low density: extended atmosphere



Detection

Flux Star + Planet Dayside Occultation Star Alone Star + Planet Nightside ➤ Time Transit Star - Planet Shadow

Image credit : Josh Winn

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Simulations



Robustness of superrotation



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Primitive equations ?

Mayne et al 2014: primitive equations hot Jupiters



Can we model warm Neptunes/super Earths with the primitive equations ?



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Mayne et al. 2019: what about warm Neptunes ?

Primitive equations

Much faster computationnally than full equations.

Four approximations, based on the Earth's shallow atmosphere/oceans:

- 1) Hydrostatic balance. Never a problem in our simulations.
- 2) Shallow fluid approximation: $r \approx R$, $\partial/\partial r \approx \partial/\partial z$
- 3) Traditional approximation: buyoancy dominates Coriolis. No latitudinal component of rotation
- 4) Gravity is constant with height.



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Assumptions

Goal: estimate the validity of $w \ll v \tan(\varphi)$

Analytical estimates based on four assumptions:

- 1) The atmosphere is globally superrotating
- **2)** $\Delta T \ll \Delta T \downarrow forcing$

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- 3) V is only due to Coriolis
- 4) Incompressible hydrostatic atmosphere



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0.3 bars

II. Warm Neptunes Results

$$\begin{split} U \sim \sqrt{3\&2} \ \pi \ R \downarrow p \ \mathscr{R} \Delta T \downarrow forcing / \tau \downarrow rad & \sim 1400 \ m.s \uparrow -1 \\ W \sim H/L \ \sqrt{3\&2} \ \pi \ R \downarrow p \ \mathscr{R} \Delta T \downarrow forcing / \tau \downarrow rad \\ \tan(\varphi) \sim \pi/2 \ R \downarrow p \ \Omega \sin \uparrow 2 \ (\varphi) / \cos(\varphi) \end{split}$$

Increases with: radius, forcing Decreases with; molecular weight Increases with: radius, rotation rate

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V

No dependence on gravity



P



-30.0000 -20.0000 -10.0000 0.00 666.667 1333.33 2000.00



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Verification





0.00 666.667 1333.33 2000.00

1000.00



(c) dT+ Full: 800-1000 days

Verification



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300

350

350

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Conclusions

Primitive equations: large planet moderately forced with heavy molecular weight

Counter intuitive: large rotation rate.

Impact: phase curves.

In the limit of global superrotation.

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Maximum temperature



Corot 2b

The inflated radius

Simple evolution models (Guillot et al 1996 + Goukenleuque et al 2000): HD209458b bigger than expected

Since then:



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Showman Polvani 2011

Solar system: superrotation sometimes associated with propagation and dissipation of Rossby waves

Hot Jupiters: Rossby deformation radius ~ planet size



Beta plane shallow water

 $\partial u/\partial t - yv + \partial h/\partial x + u/\tau \downarrow drag = 0,$

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 $\partial v/\partial t + yu + \partial h/\partial y + v/\tau \downarrow drag = 0,$

 $\partial h/\partial t + \partial u/\partial x + \partial v/\partial y + h/\tau \downarrow rad = Q.$

Timeline from SP11



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Superrotation in hot Jupiters

Equilibration of the jet: sequence of linear steady states, Tsai et al. 2014. Vertical tilt of the wave



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Limits of SP11 and Tsai et al. 2014

Tsai: Need for an initial superrotation followed by slow evolution. Given by SP11 ?

Linear steady state of SP11 with appropriate $\tau \downarrow drag$ and $\tau \downarrow rad$: (komacek & showman 2016)

Linear steady state never reached

Need for time dependent considerations



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Debras et al.2019, accepted

Linear time dependent solution:

 $X \downarrow F = \sum n \uparrow [] q \downarrow n X \downarrow n / \sigma \downarrow n - i \omega \downarrow n (1 - e \uparrow (i \omega \downarrow n - \sigma \downarrow n)t)$

Steady state dominated by Rossby waves

BUT

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Limit of short timescales: Rossby and Kelvin waves with comparable amplitudes $X \downarrow F \approx \sum n \uparrow [] q \downarrow n X \downarrow n t$



Rossby and **Kelvin** waves

Rossby: ~ zero pressure at the equator, rotating winds



Kelvin: no meridional winds, maximum pressure at the equator

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Intermediate and steady linear state

Steady: Rossby dominated

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Short timescales state: mix R-K, Similar to SP11 steady state

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More accurate timeline



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Quasi-linear studies and accelerations

Separation of linear and non linear considerations too simple

Quasi linear/statistical studies ? Srinivasan & Young 2012, Bouchet et al. 2013, Bakas et al. 2015

Vertical accelerations:

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Summary



Can we model warm Neptunes/super Earths with the primitive equations ?

Range of planets where traditional approximation breaks

Analytical priors verified numerically

Strong impact on the comparison with observations

What is the physical origin of equatorial superrotation on hot Jupiters ?

Initial phases of simulated superrotation not perfectly understood

Time dependent linear processes needed to be taken into account

Crucial link between superrotation and interior profiles

Thank you !





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Perspectives



Models compatible with Saturn ?

Love numbers with CMS method

Interior of hot Jupiters: flatter temperature gradient, favorable for semi convection

Radius re-inflation for a non convective planet?





Moore et al. 2018

Post doctoral position in Toulouse

EOS





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Why is gravity interesting ?



$$J_{2k} = -\frac{4\pi}{MR^{2k}} \int_0^R \int_0^1 \rho(\vec{r'}) r'^{2k+2} P_{2k}(\mu') d\mu' dr'$$

Pioneer 10-11: 1973-1974 Voyager 1-2: 1978-1979 Galileo: 1995

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Diluted Core

Moll et al. 2017



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III. New models of the interior of Jupiter



Immiscibility

Metallic H/He immiscibility possible

Helium rain: heating up the planet Stevenson Salpeter 1977

Decrease in Z < 10%



II. Superrotation

Solar system

Mid latitudes jet (Jupiter, Saturn, Earth) : Rossby waves

Global superrotation (Venus) : Meridional momentum diffusion

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Equatorial jet : Rossby waves from magnetic field of heat fluxes

II. Spin-up of superrotation

Tsai et al. 2014

3D = infinite sum of 2D with different equivalent depths (wu et al. 2000) : projection of the heating function in the vertical



Equilibration of the jet from the vertical structure

II. Spin-up of superrotation

My contribution : time dependent solution

$$\begin{split} \frac{\partial u}{\partial t} &-yv + \frac{\partial h}{\partial x} + \frac{u}{\tau_{\text{drag}}} = 0, \\ \frac{\partial v}{\partial t} &+yu + \frac{\partial h}{\partial y} + \frac{v}{\tau_{\text{drag}}} = 0, \\ \frac{\partial h}{\partial t} &+ \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{h}{\tau_{\text{rad}}} = Q, \end{split} \qquad \begin{aligned} (\mathrm{i}\omega + \frac{1}{\tau_{\text{drag}}})u - yv + \mathrm{i}kh = 0, \\ (\mathrm{i}\omega + \frac{1}{\tau_{\text{drag}}})v + yu + \frac{\partial h}{\partial y} = 0, \\ (\mathrm{i}\omega + \frac{1}{\tau_{\text{rad}}})h + \mathrm{i}ku + \frac{\partial v}{\partial y} = 0. \end{aligned}$$

Homogeneous solution :

$$X_{\rm H} = \sum_{n,l} \alpha_{n,l} X_{n,l}, \quad \text{With} \quad X_{n,l} = \tilde{X}_{n,l} (x,y) e^{(i\omega_{n,l} - \sigma_{n,l})t}$$

Forced solution :

$$X_{\rm F} = \sum_{n,l} \frac{q_{n,l} \tilde{X}_{n,l}}{\sigma_{n,l} - i\omega_{n,l}} \left(1 - e^{(i\omega_{n,l} - \sigma_{n,l})(t)} \right)$$

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II. Spin-up of superrotation

Non linear considerations

$$u_{
m max} \sim rac{L}{ au_{
m drag}}$$

For real forcing, umax >> usteady : linear steady state never reached

Timescale analysis : limit of short times, Kelvin and Rossby waves have same amplitude different from limit of long times (Rossby dominate)



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