

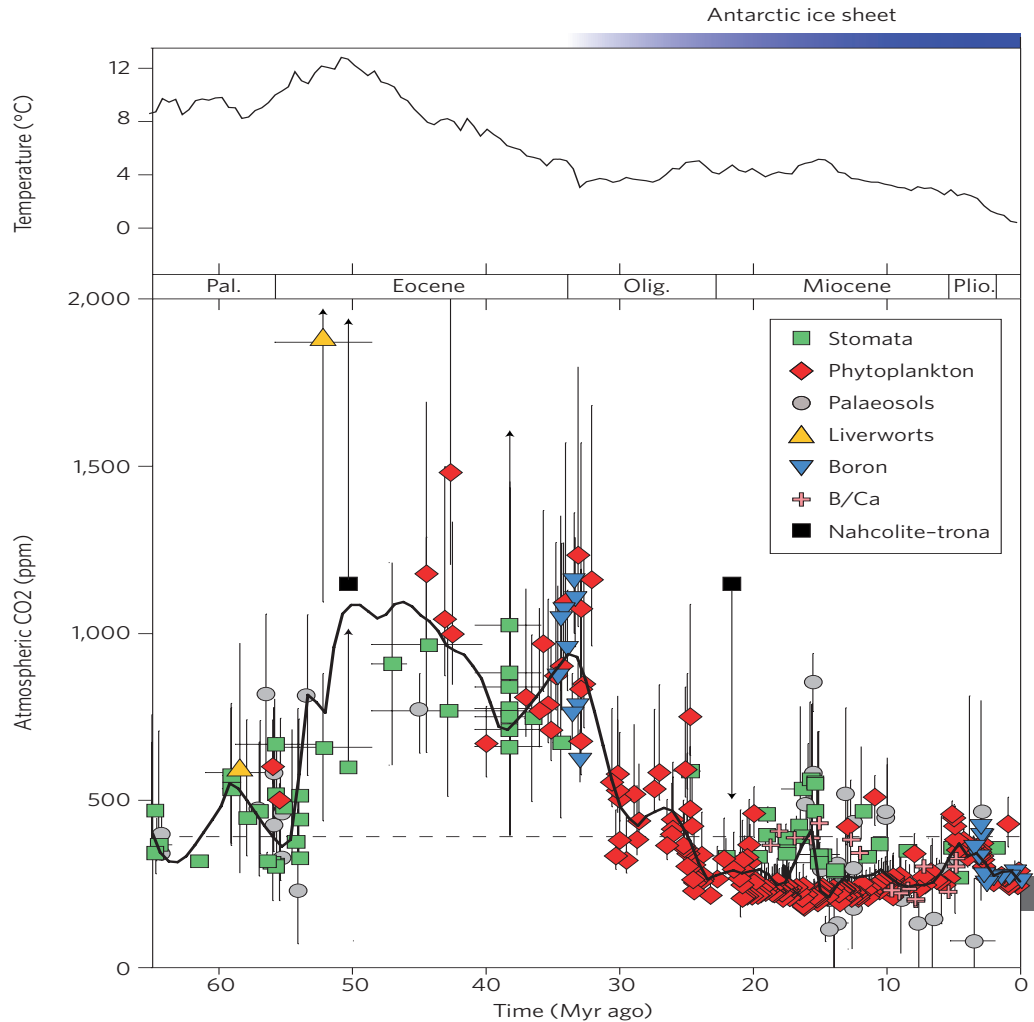
# Superrotation at Earth's surface

Rodrigo Caballero, Henrik Carlson

Department of Meteorology, Stockholm University



# The past 50 million years on Earth



Temperature at the bottom of the ocean

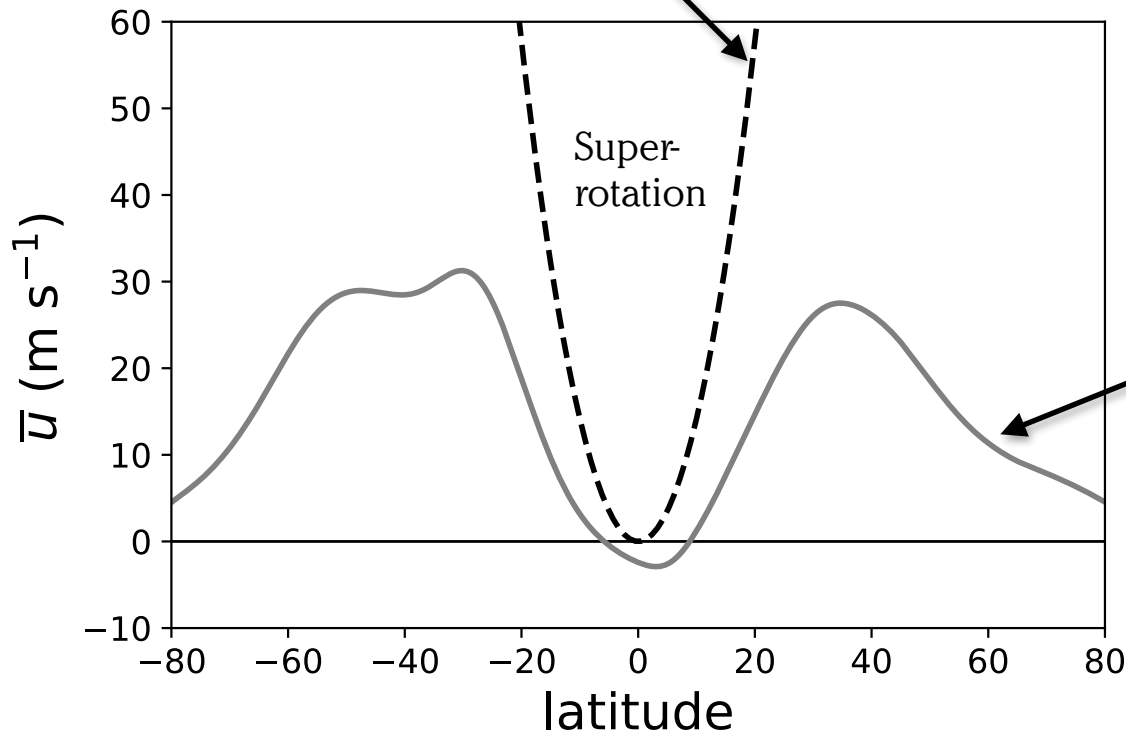
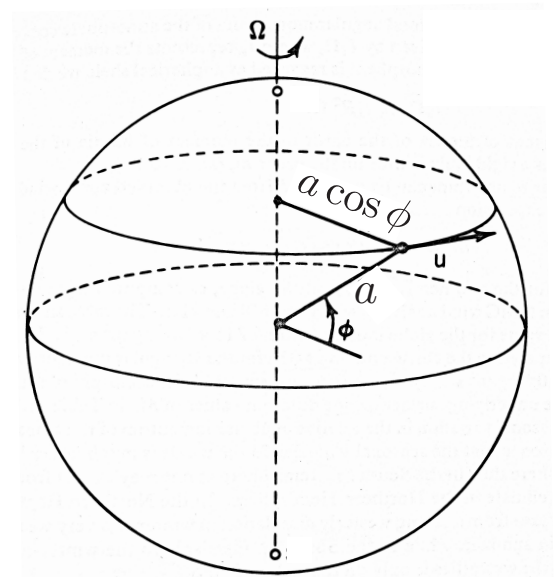
Atmospheric CO<sub>2</sub>

# Superrotation

A state in which the atmosphere has greater axial angular momentum than equatorial solid-body rotation:

$$M = \Omega a^2 \cos^2 \varphi + \bar{u} a \cos \varphi > \Omega a^2$$

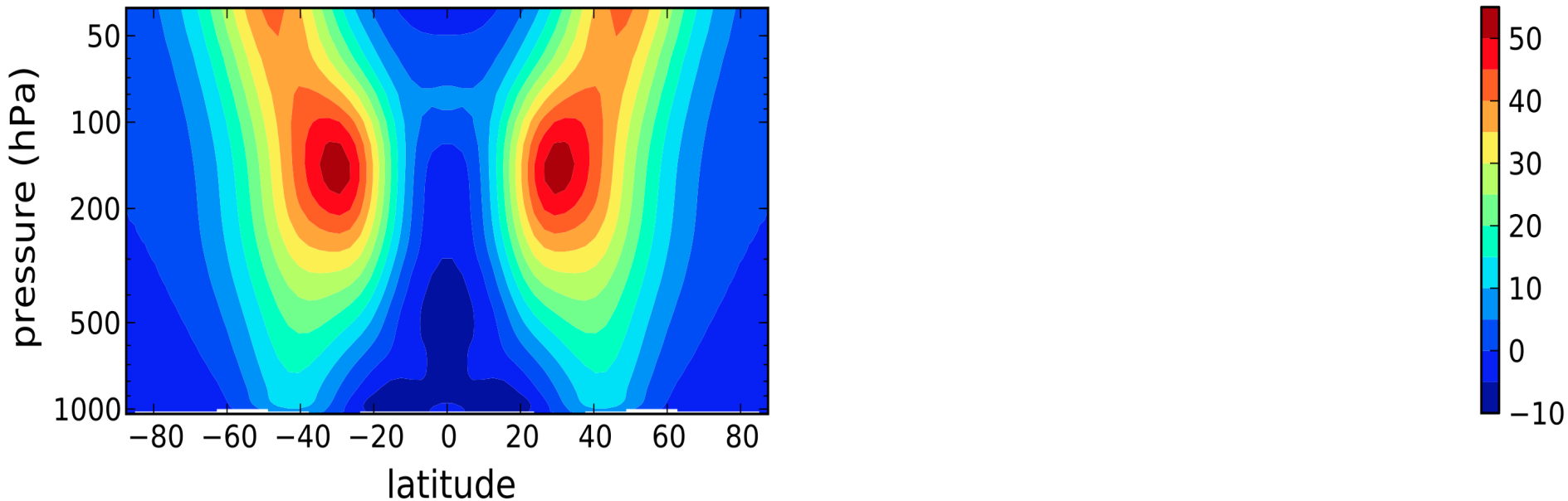
$$\bar{u} > \Omega a \frac{\sin^2 \varphi}{\cos \varphi}$$



# Superrotation in warm climates

Zonal-mean zonal wind in simulations with CAM3 coupled to aquaplanet slab ocean

COLD (1xCO<sub>2</sub>)



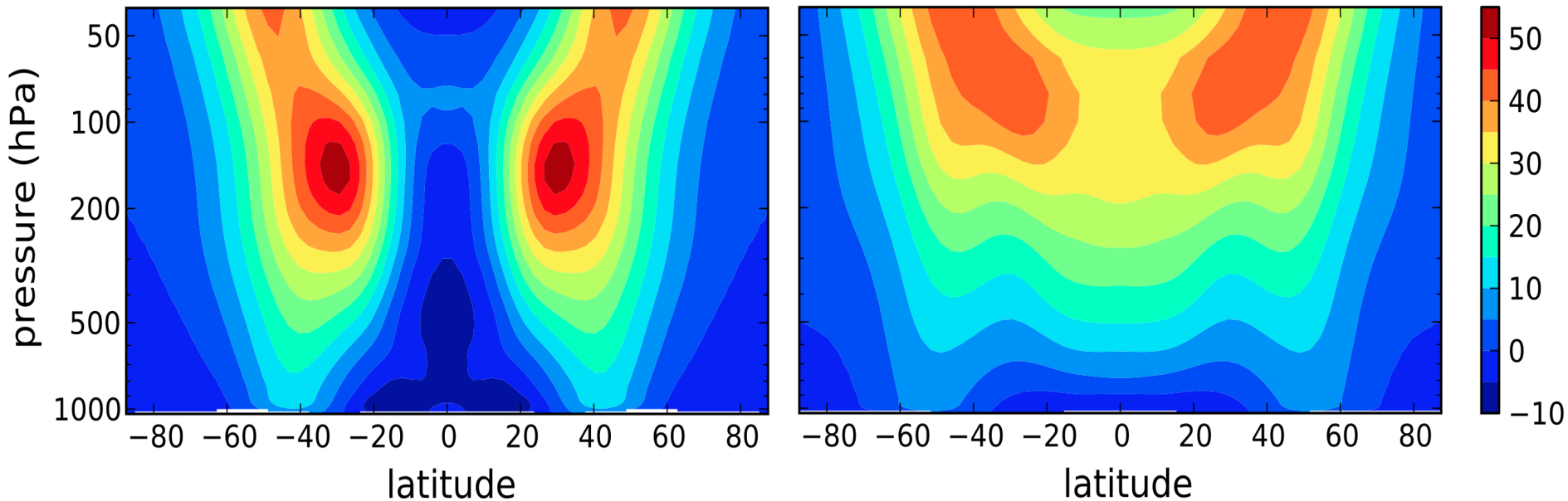


# Superrotation in warm climates

Zonal-mean zonal wind in simulations with CAM3 coupled to aquaplanet slab ocean

COLD ( $1\times\text{CO}_2$ )

HOT ( $32\times\text{CO}_2$ )



(Caballero & Huber GRL 2010)

# What drives superrotation?

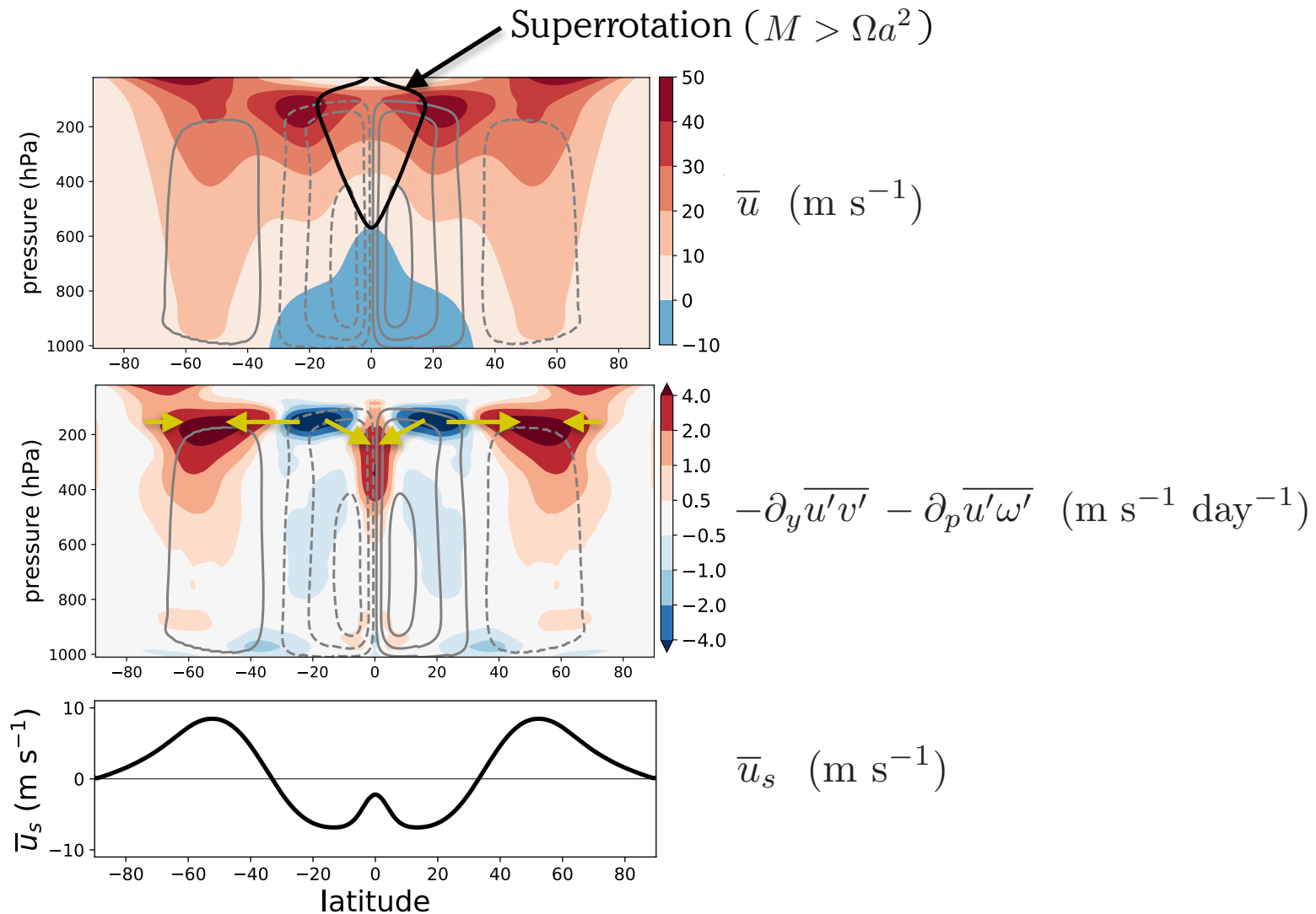
Zonal-mean momentum budget:

$$\partial_t \bar{u} = \underbrace{(f - \partial_y \bar{u}) \bar{v} - \bar{\omega} \partial_p \bar{u}}_{\substack{\text{mean flow} \\ \text{(conserves } M)}} - \underbrace{\partial_y \overline{u'v'} - \partial_p \overline{u'\omega'}}_{\text{large-scale eddies}} - \underbrace{g \partial_p \bar{\tau}}_{\text{viscosity}}$$

- ▶ Mean flow conserves angular momentum
- ▶ Downgradient (diffusive) eddy momentum fluxes can only dilute  $M$  maxima
- ▶ Superrotation requires *countergradient* eddy momentum fluxes by large-scale eddies

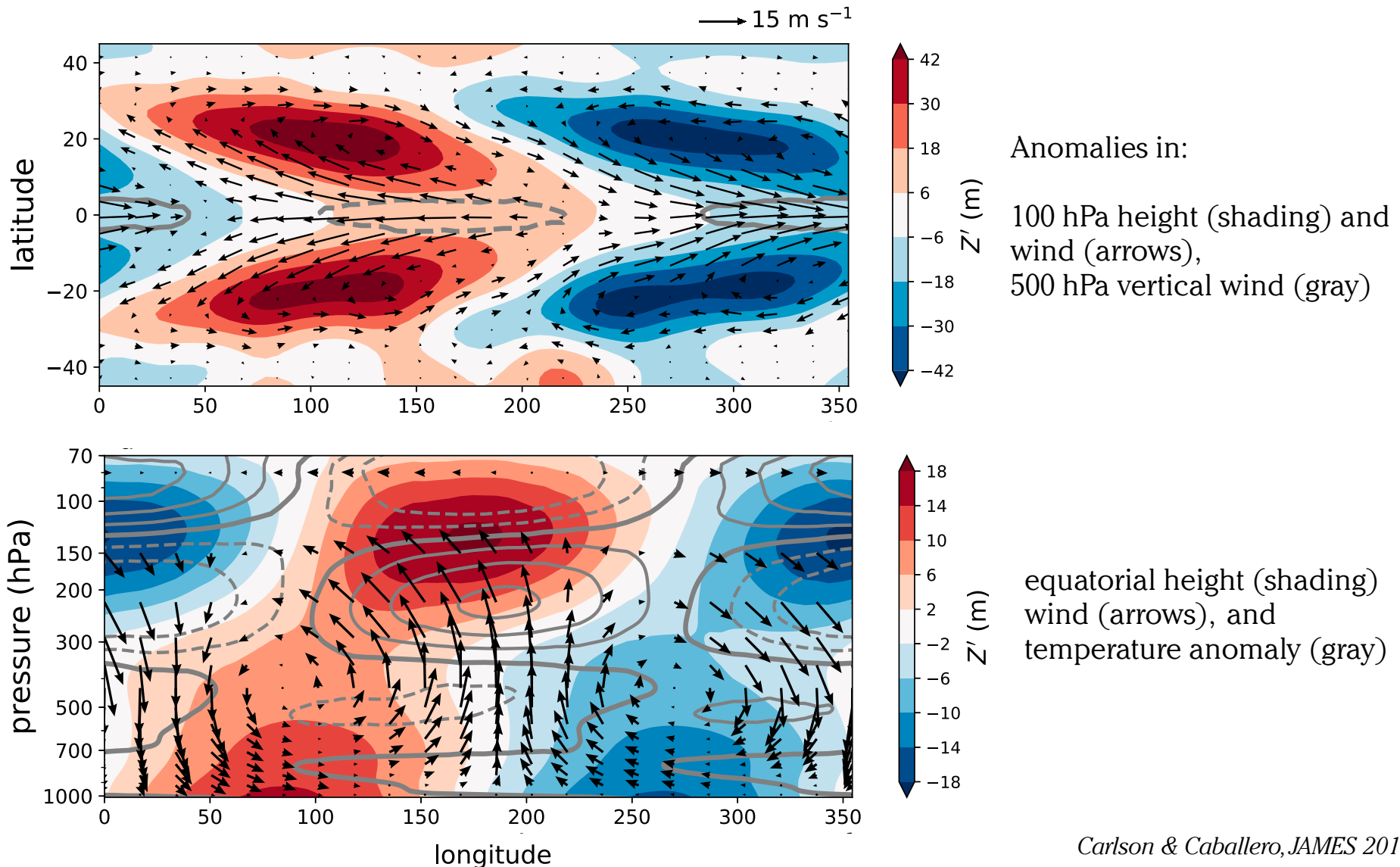
# Tropical eddies drive warm-climate superrotation

Aquaplanet (CAM4) with very warm surface temperature

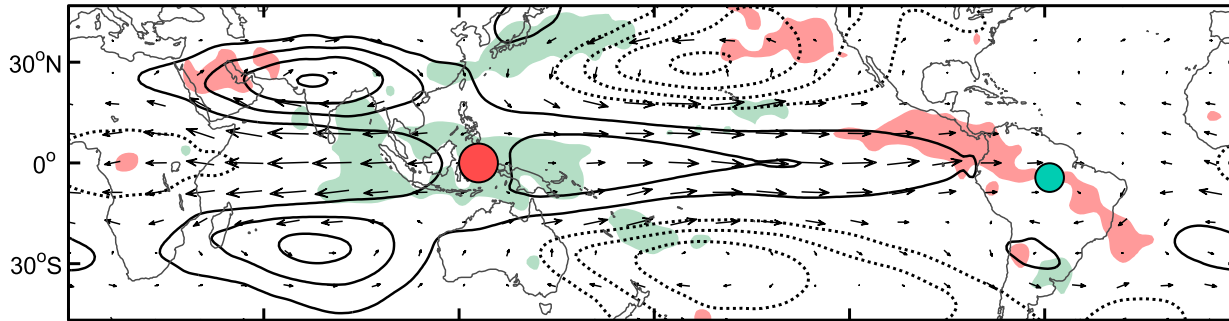


# What are these tropical eddies?

- ▶ Leading EOF of tropical variability looks a lot like the Madden-Julian Oscillation

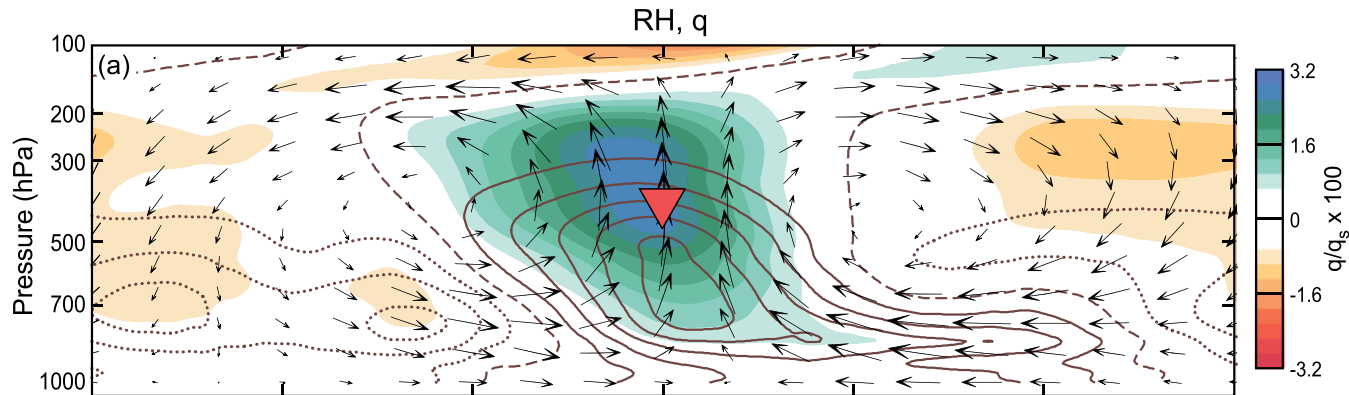


# Observed MJO



Anomalies in:

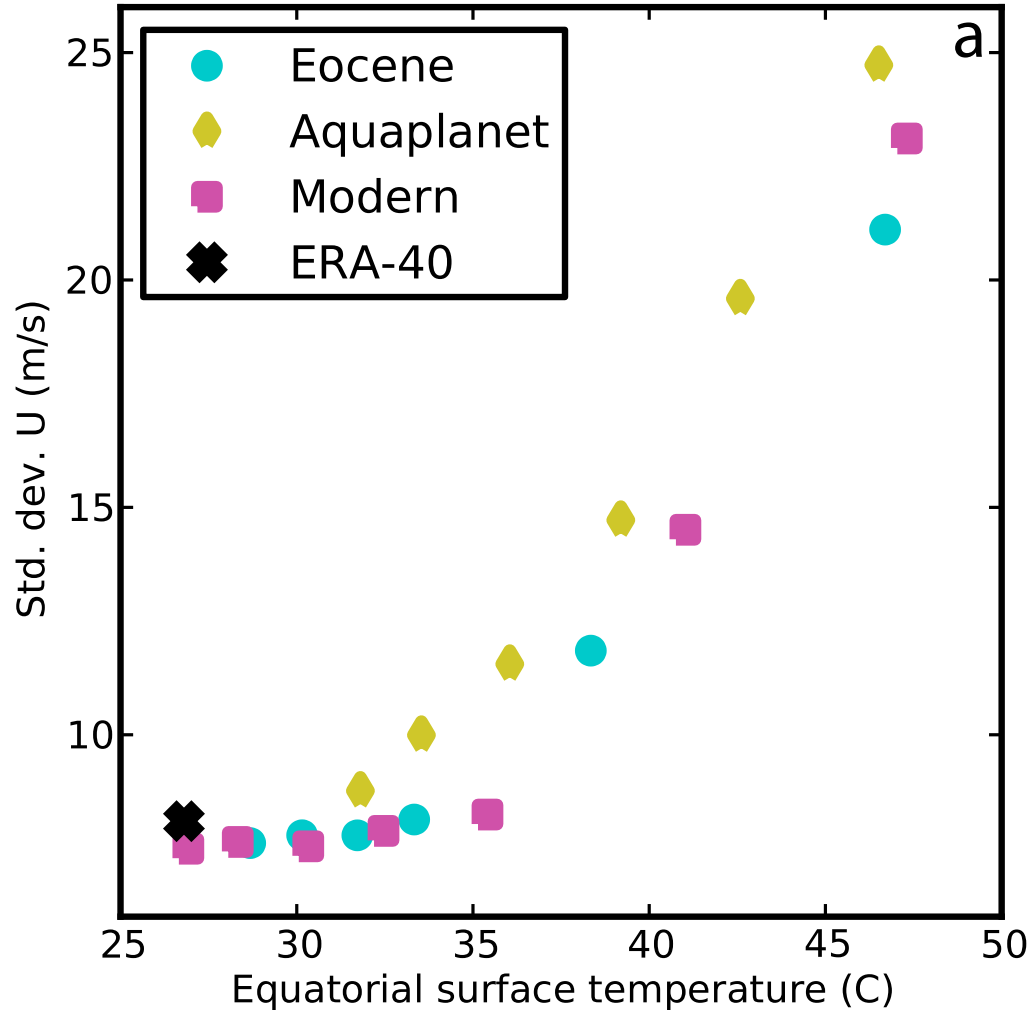
upper troposphere  
height (contours) and  
wind (arrows),  
400 hPa vertical wind (color)



equatorial wind (arrows)  
and humidity (shading)

# Enhanced MJO in warm climates

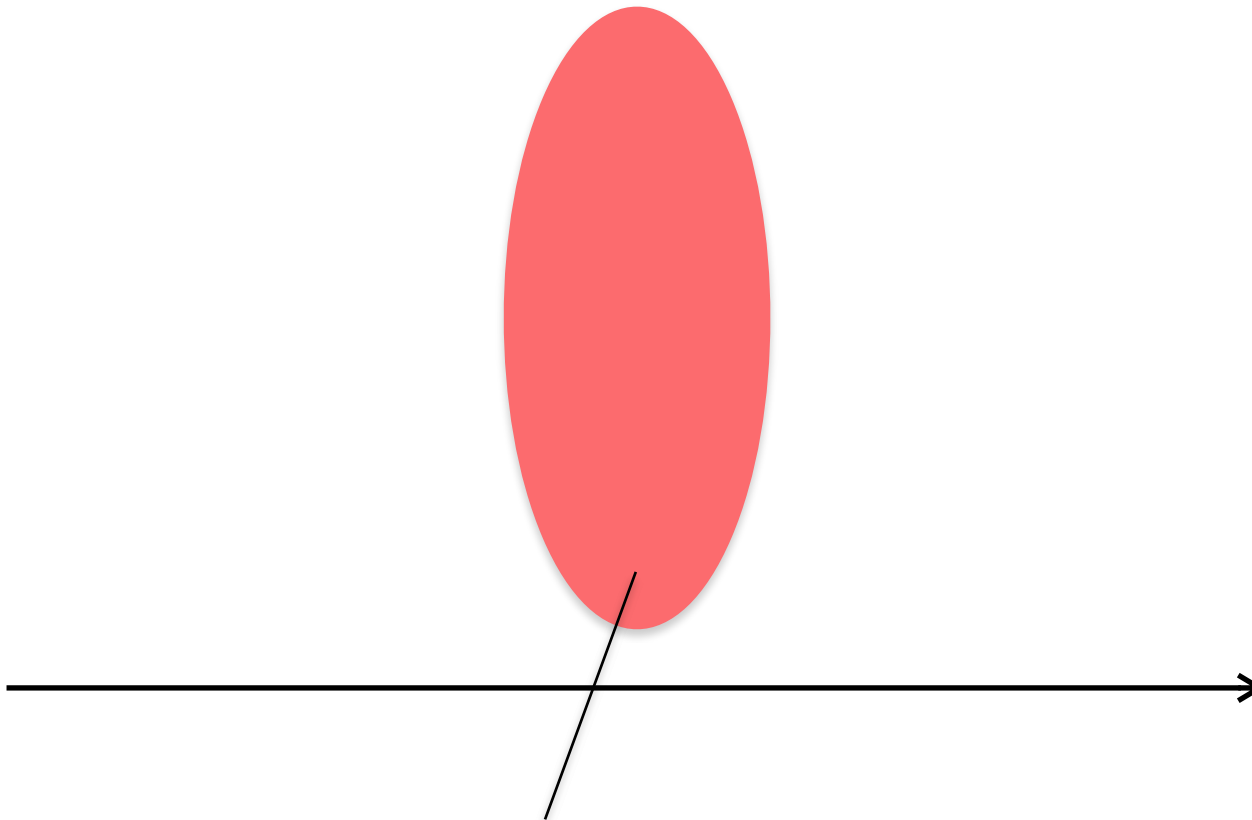
Full-complexity GCM simulations (CAM3+slab ocean) with increasing CO<sub>2</sub>



## 2 topics for the rest of this talk:

- ▶ Why does MJO amplitude increase in warm climates?
- ▶ Surface superrotation: Can equatorial winds be westerly at the surface?

# “Moisture-mode” instability

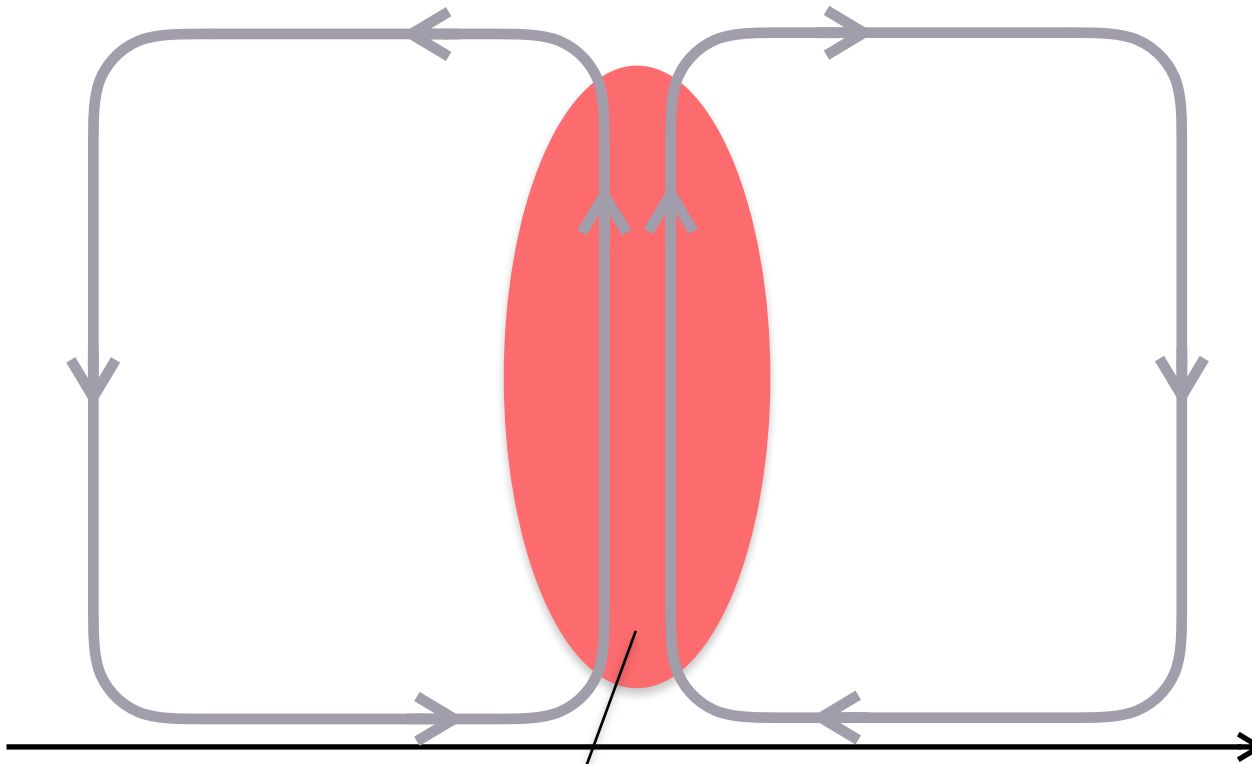


Moist static energy anomaly

$$h = C_p T + gz + Lq$$



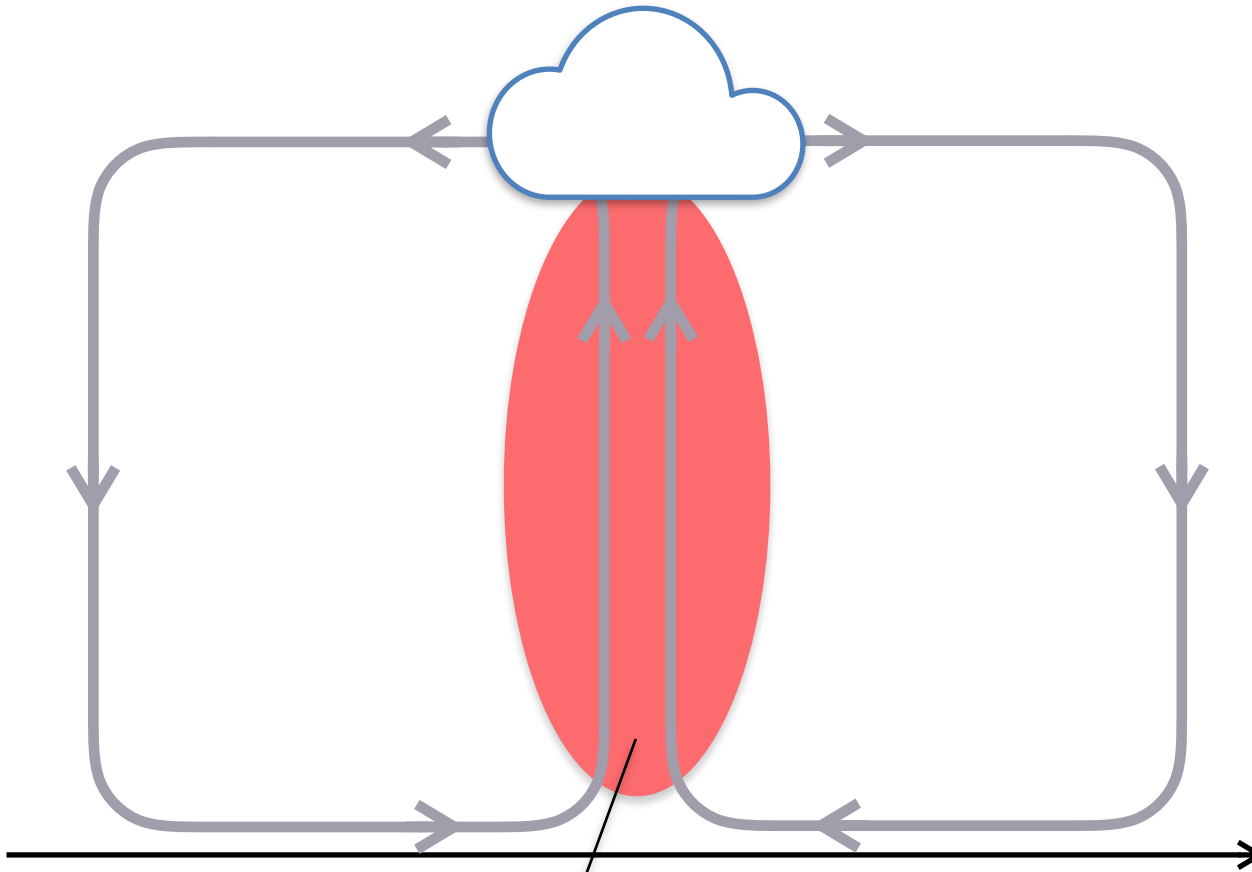
# “Moisture-mode” instability



Moist static energy anomaly

$$h = C_p T + gz + Lq$$

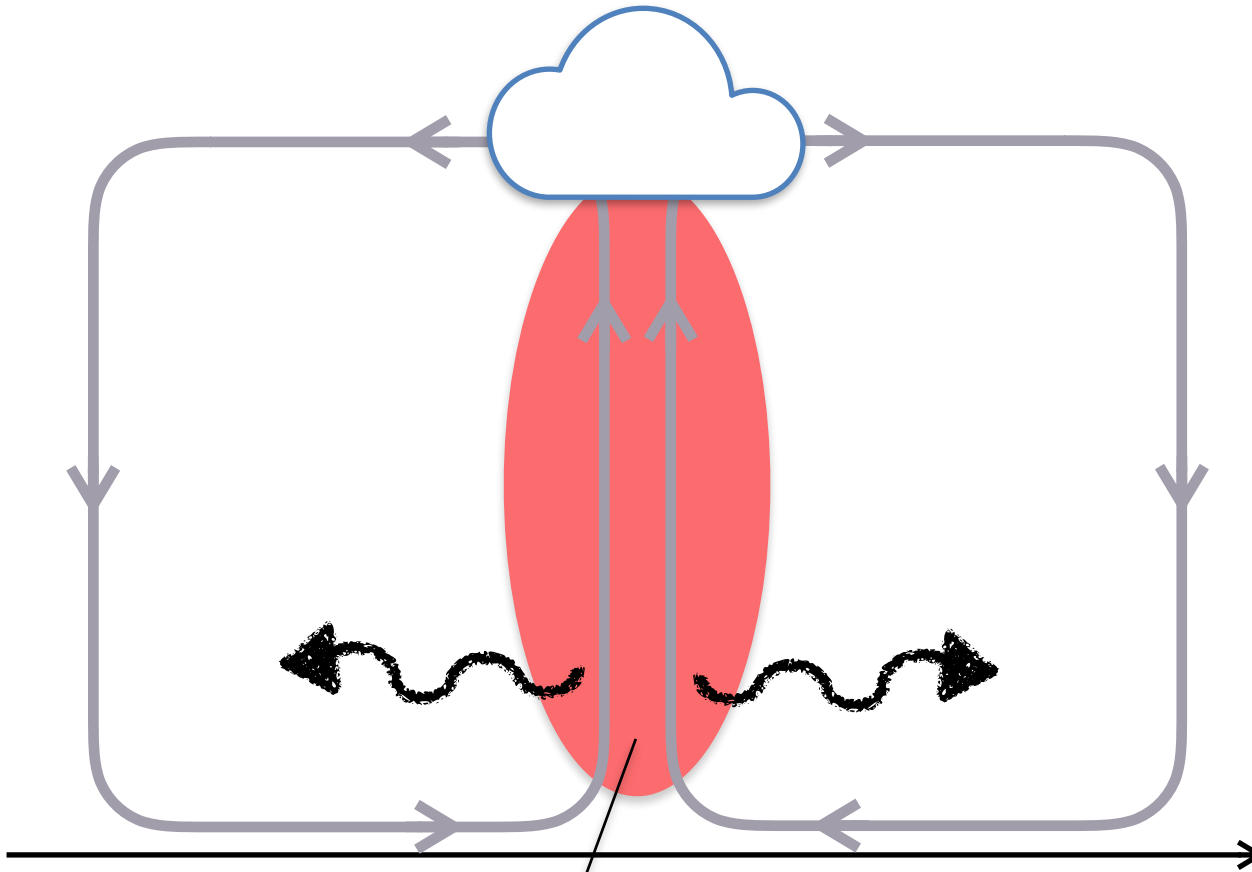
# “Moisture-mode” instability



Moist static energy anomaly

$$h = C_p T + gz + Lq$$

# “Moisture-mode” instability



Moist static energy anomaly

$$h = C_p T + gz + Lq$$

# Moist static energy budget

Local evolution of moist static energy  $h = C_p T + gz + Lq$

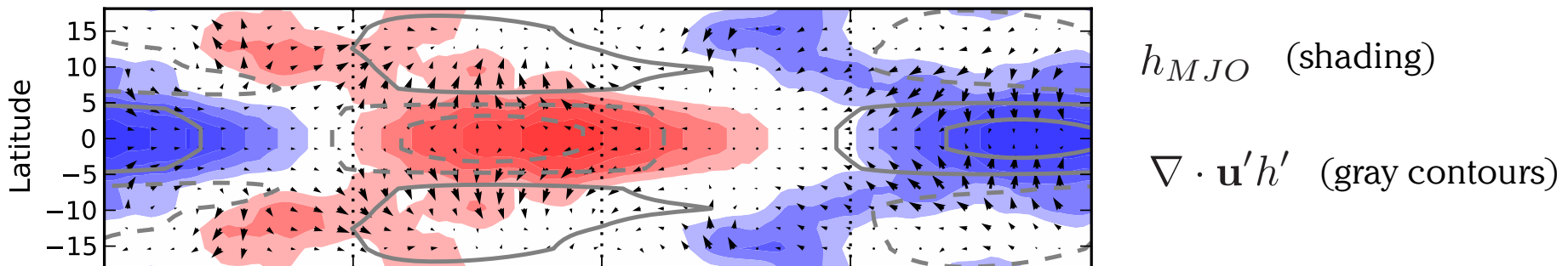
$$\partial_t h = -\mathbf{u} \cdot \nabla h - \omega \partial_p h + R + F_s$$

Partition into zonal mean, MJO mode and residual:

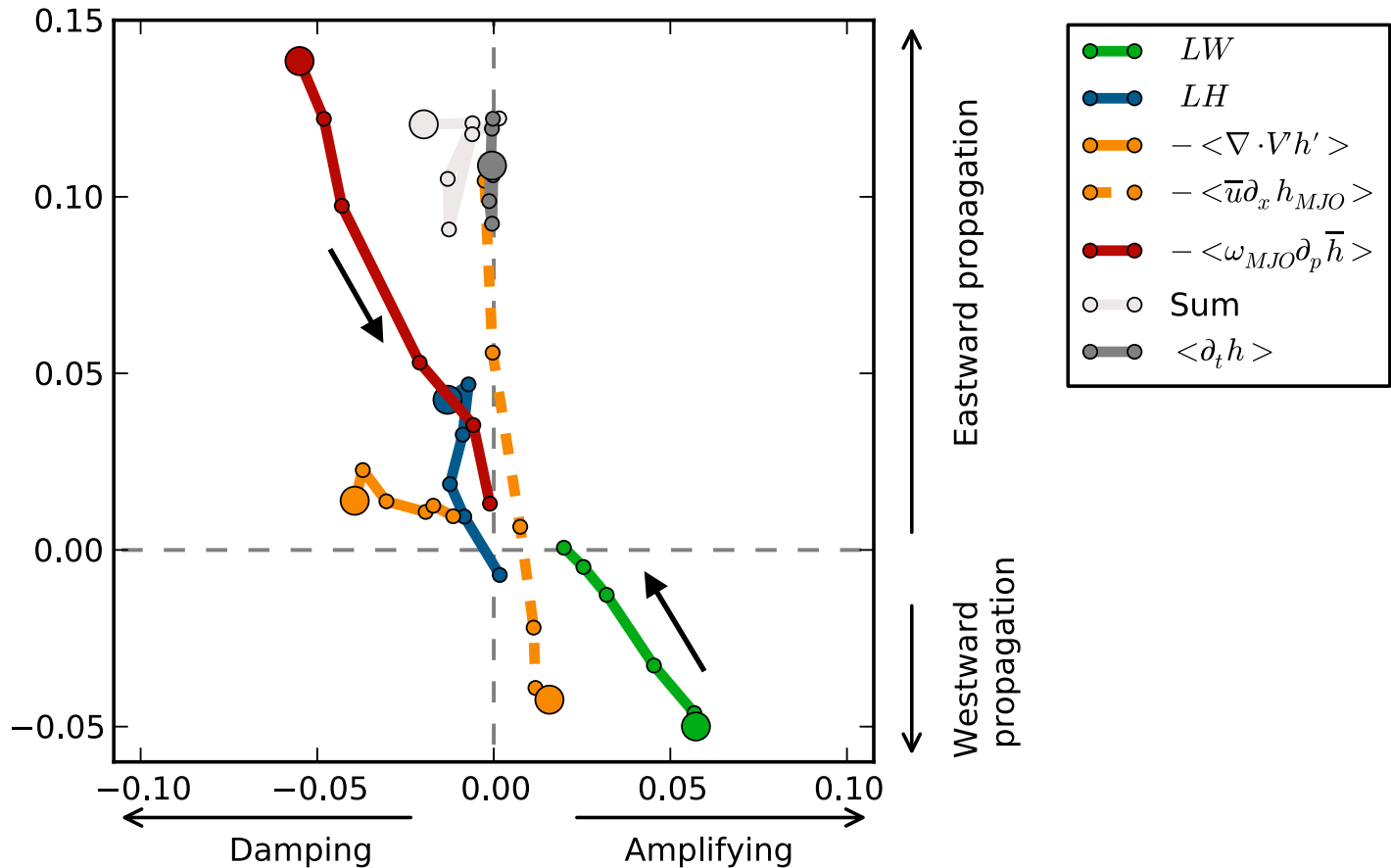
$$h = \bar{h} + h_{MJO} + h'$$

Substitute into evolution equation, project onto MJO mode; turns out that

$$\partial_t h_{MJO} \approx \underbrace{-\omega_{MJO} \partial_p \bar{h}}_{\text{Negative (damping)}} - \underbrace{\nabla \cdot \mathbf{u}' h'}_{\text{Positive (amplifying)}} + R_{MJO}$$

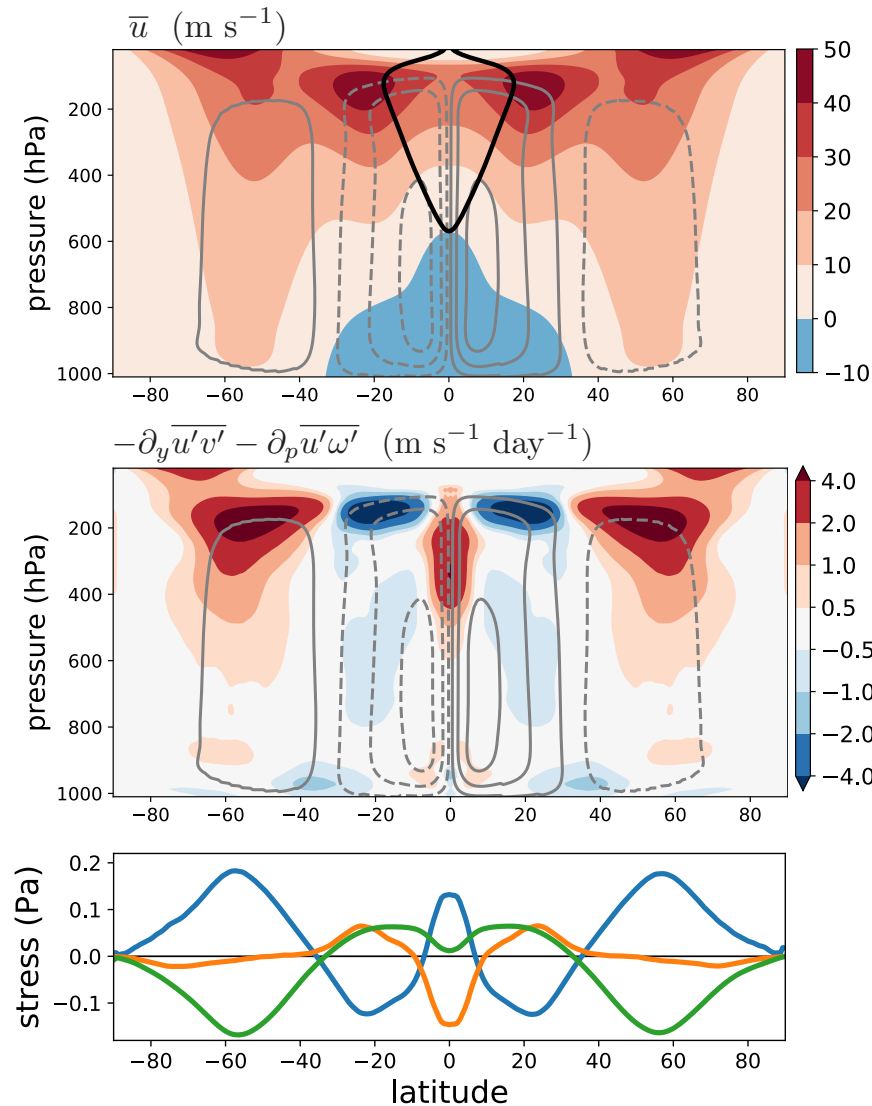


# Response of MSE budget terms to warming



Damping gets weaker as climate warms  $\Rightarrow$  MJO amplitude grows

# Momentum transport and surface wind



Momentum budget:

$$\partial_t \bar{u} = (f - \partial_y \bar{u}) \bar{v} - \bar{\omega} \partial_p \bar{u} - \partial_y \overline{u'v'} - \partial_p \overline{u'w'} - g \partial_p \bar{\tau}$$

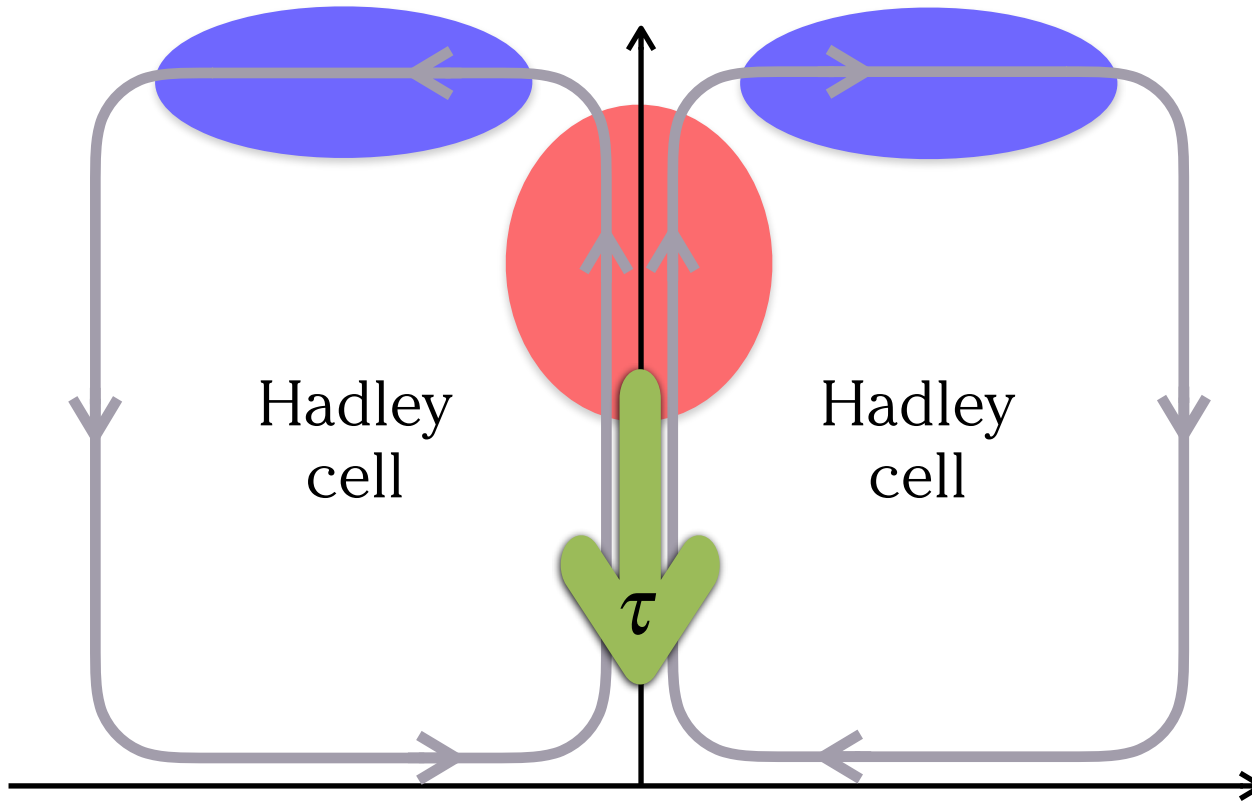
Take vertical integral  $\langle \cdot \rangle = g^{-1} \int_0^{P_s} (\cdot) dp$  :

$$-\langle \bar{v} \partial_y \bar{u} \rangle - \langle \bar{\omega} \partial_p \bar{u} \rangle - \partial_y \langle \overline{u'v'} \rangle - \bar{\tau}_s = 0$$

mean flow
large-scale eddies
surface stress

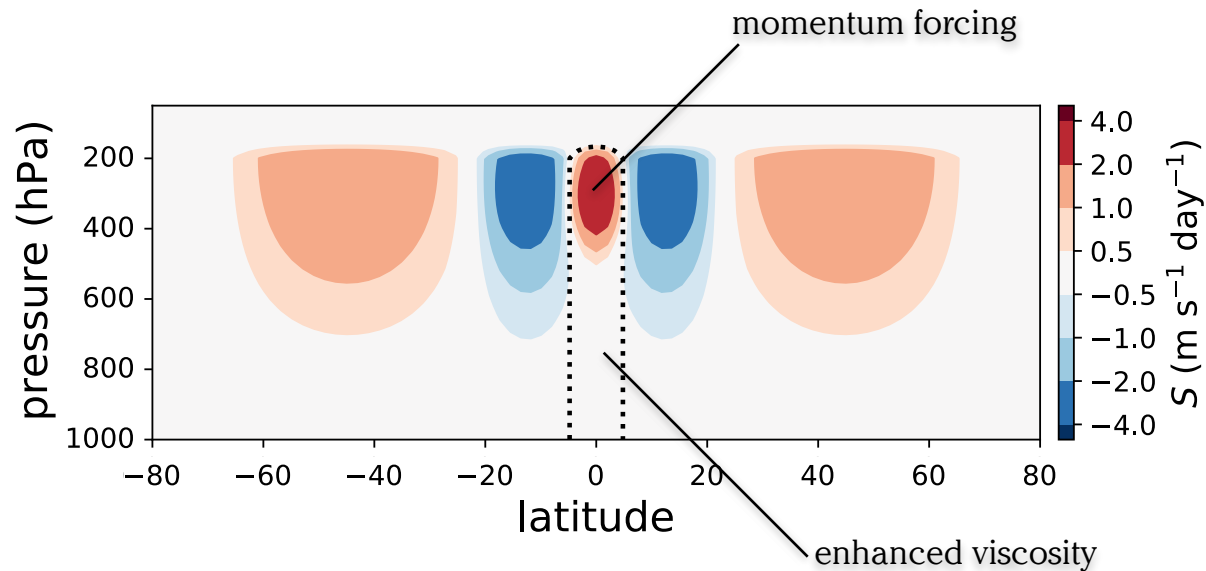
- ▶ Midlatitudes:  
eddy mom. convergence balanced by surface westerlies
- ▶ Equator:  
eddy mom. convergence balanced by surface westerlies

# A pathway to surface superrotation



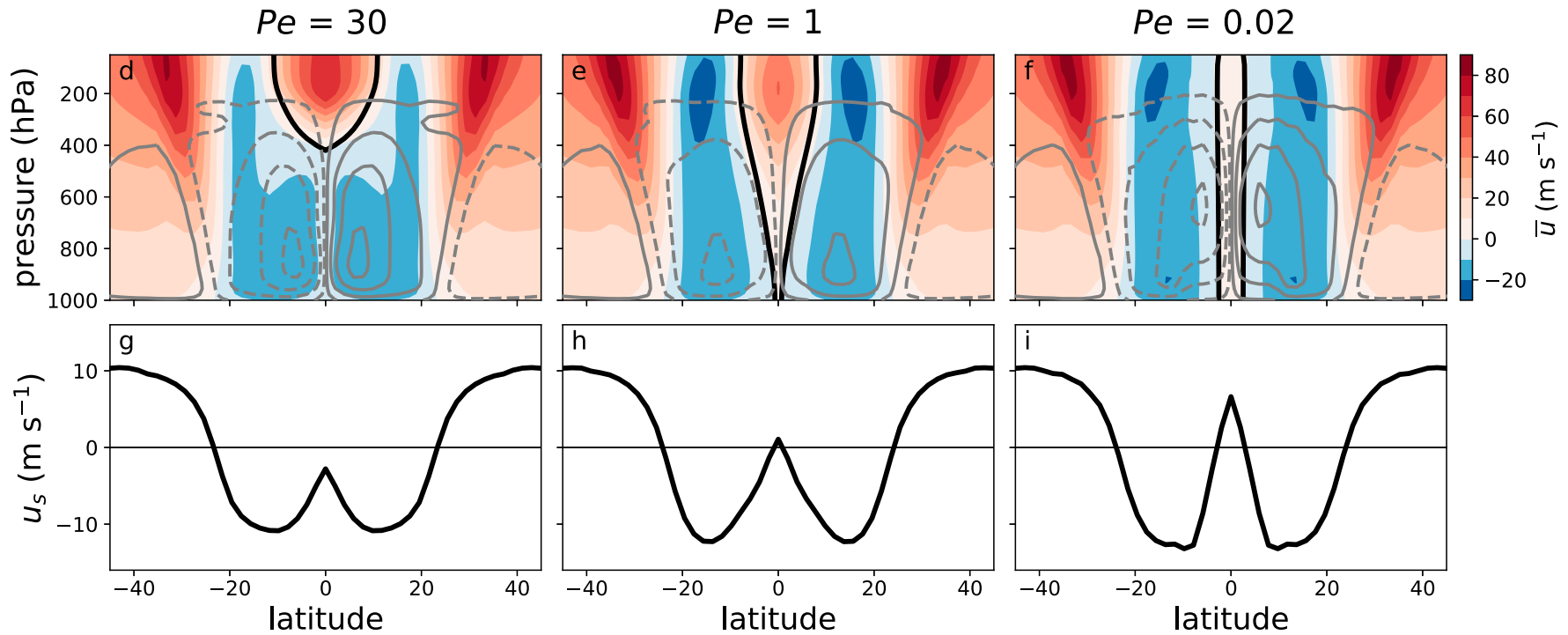
# Axisymmetric simulations

- ▶ 2D axisymmetric model, similar to Held & Hou (1980)
- ▶ Momentum forcing by prescribed distribution of zonal torque
- ▶ Diffusive vertical viscosity enhanced near equator





# Increased equatorial viscosity leads to surface superrotation



Péclet number:  $Pe = \frac{\text{advective timescale}}{\text{diffusive timescale}} = \frac{\omega/p_s}{\nu/H^2}$

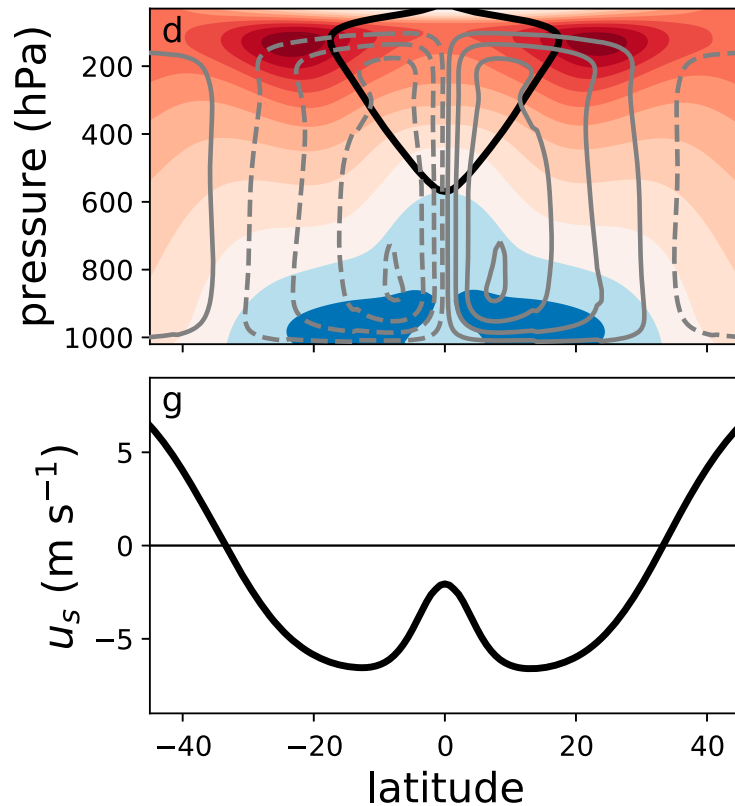
# Surface superrotation in full GCM

Reduce Péclet number in GCM by:

- ▶ reducing equator-pole temperature gradient ( $\Rightarrow$  reduced Hadley cell strength)
- ▶ artificially increasing convective momentum transport (CMT)

Control case

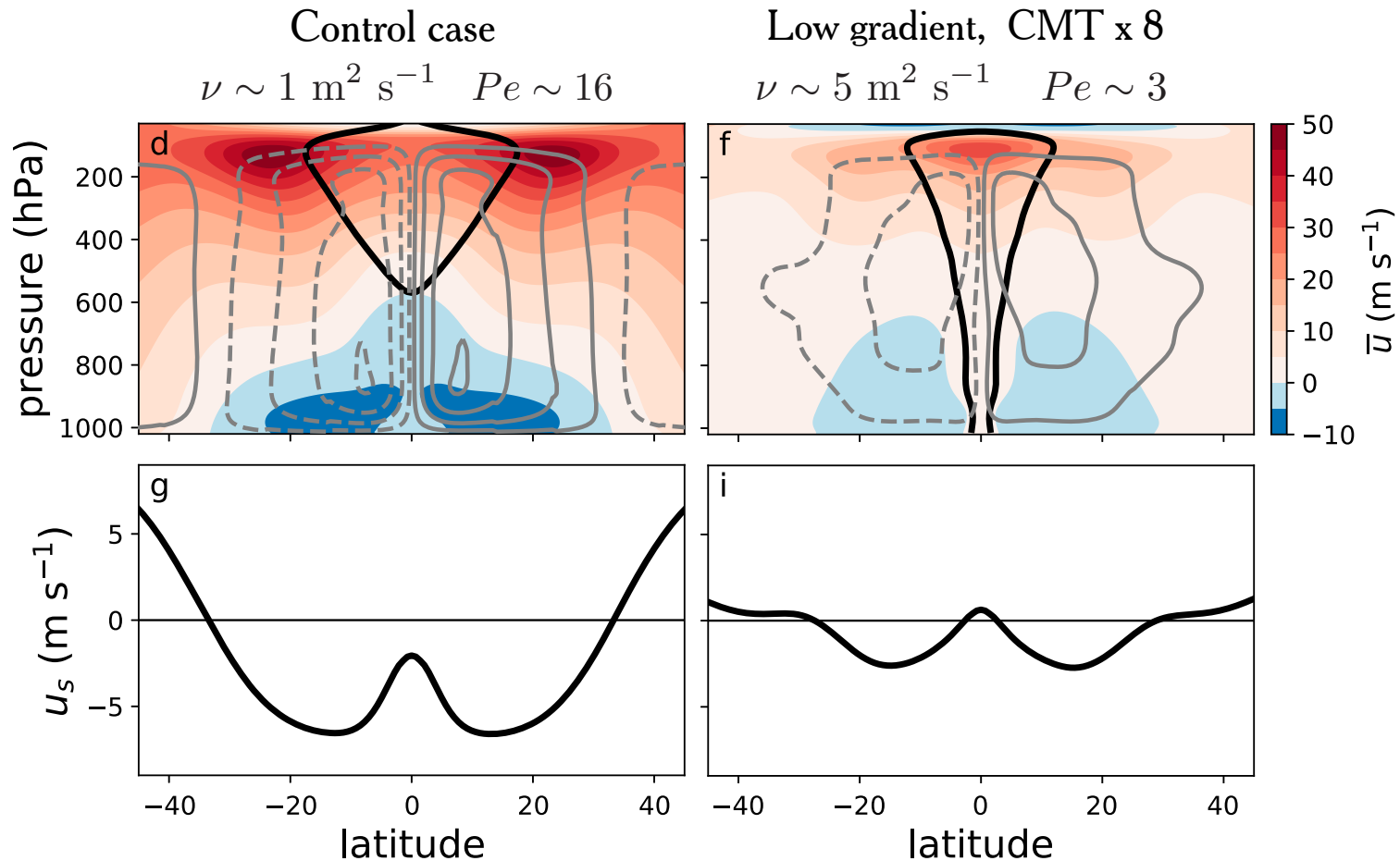
$$\nu \sim 1 \text{ m}^2 \text{ s}^{-1} \quad Pe \sim 16$$



# Surface superrotation in full GCM

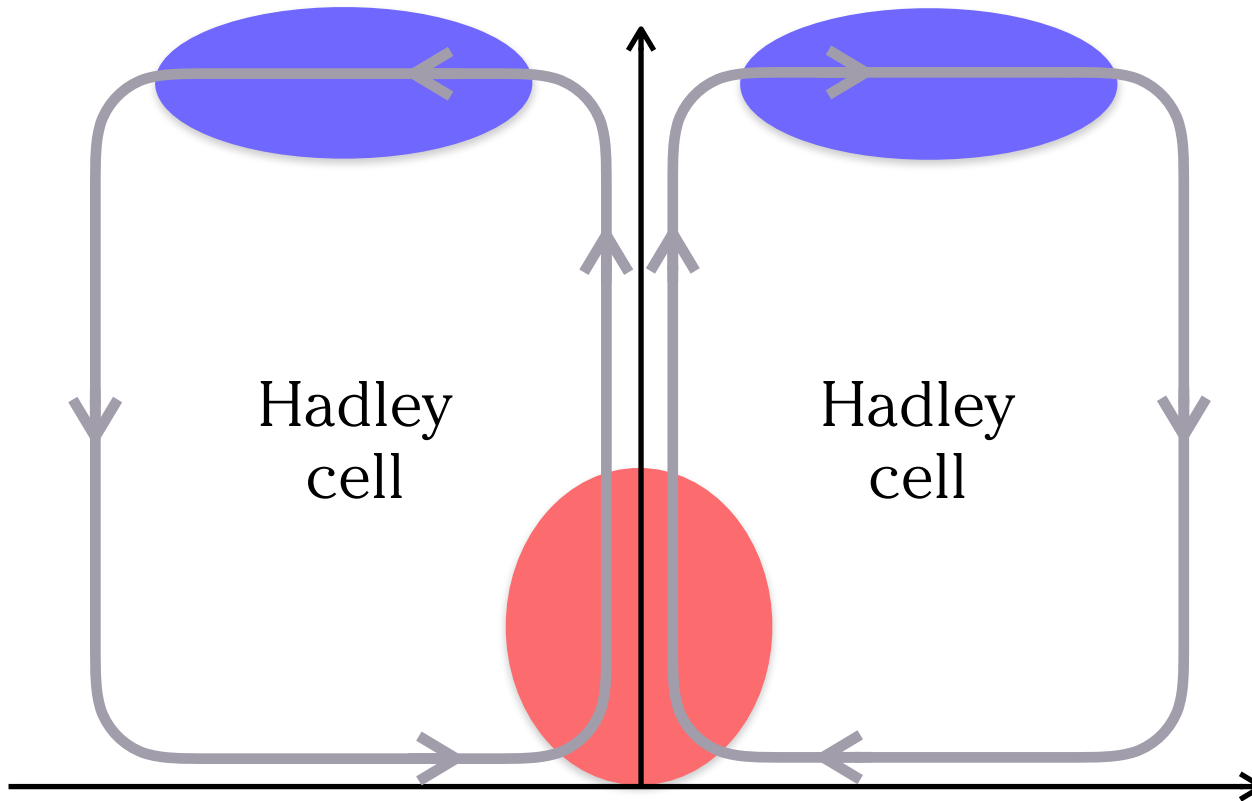
Reduce Péclet number in GCM by:

- ▶ reducing equator-pole temperature gradient ( $\Rightarrow$  reduced Hadley cell strength)
- ▶ artificially increasing convective momentum transport (CMT)



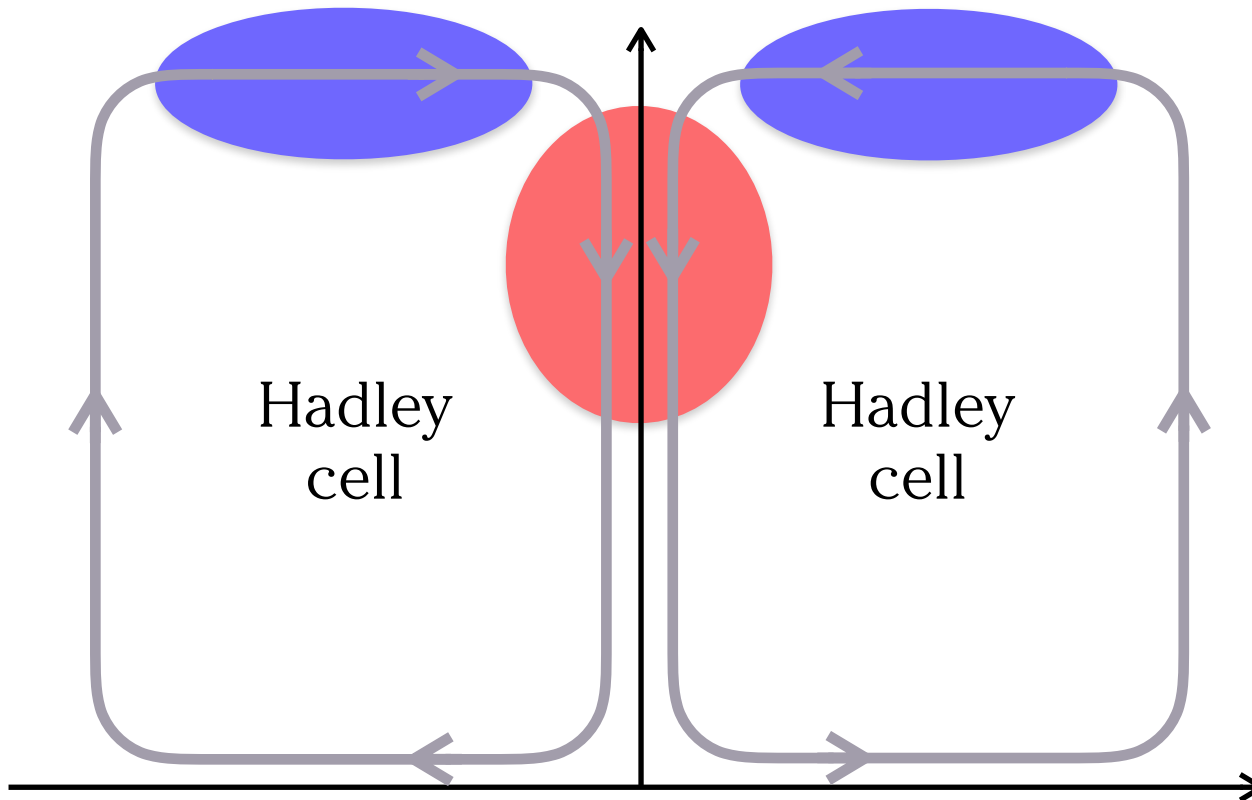
# Surface superrotation — pathway 2

Momentum convergence at surface



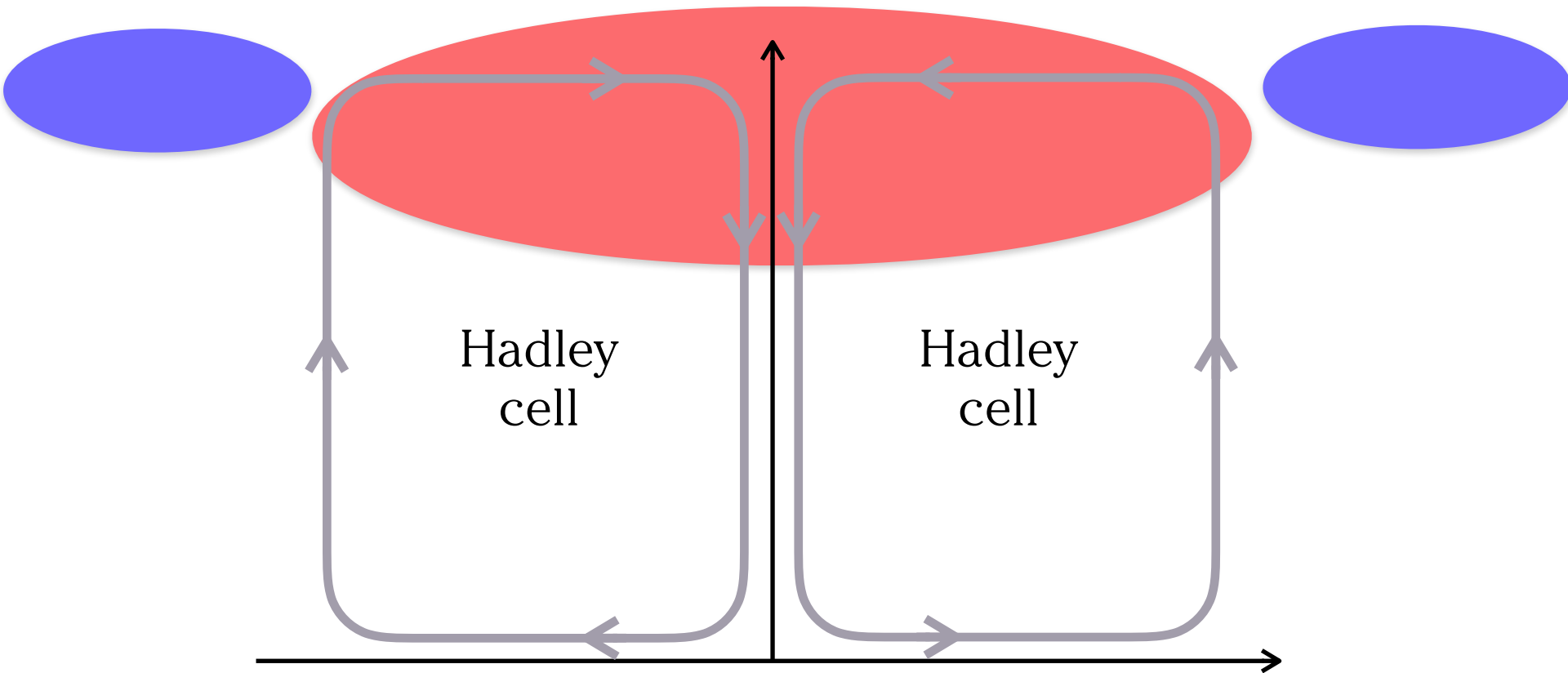
# Surface superrotation — pathway 3

Reversed Hadley cell



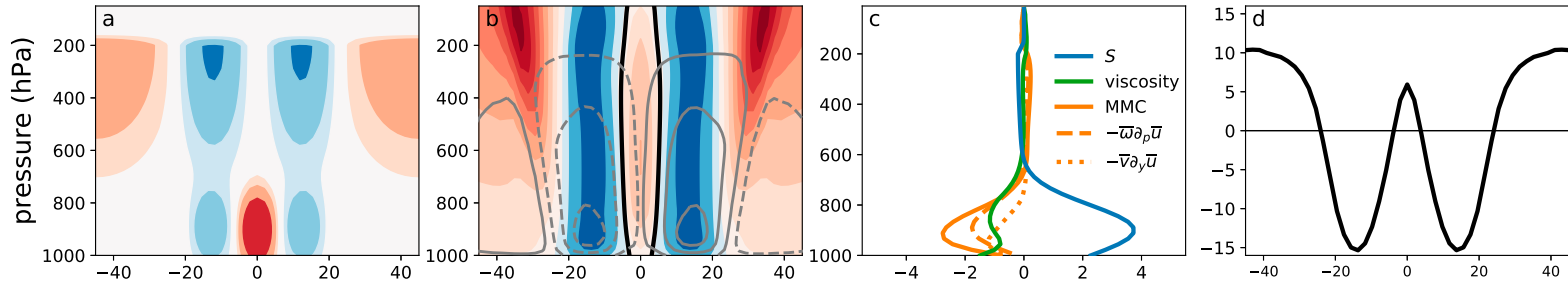
# Surface superrotation — pathway 4

Momentum convergence  
throughout Hadley cell region

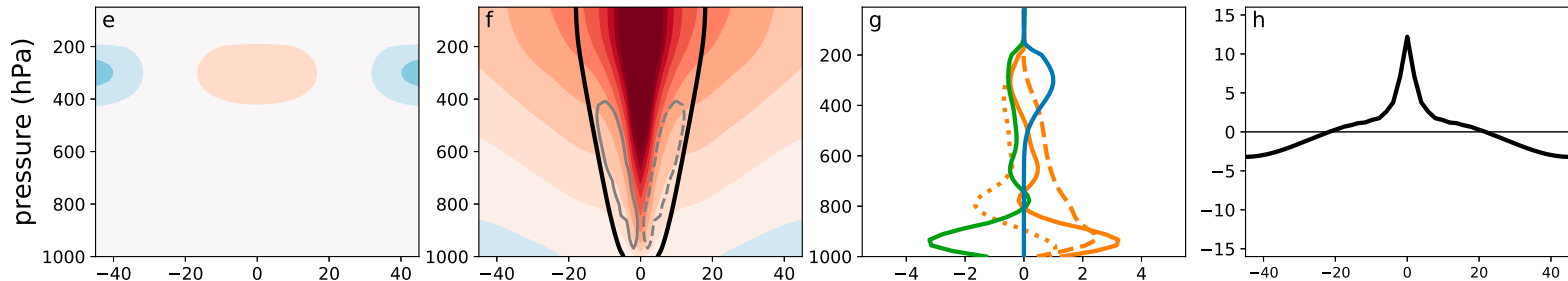


# Can realise all pathways in axisymmetric model

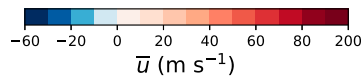
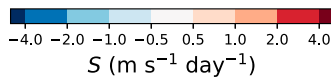
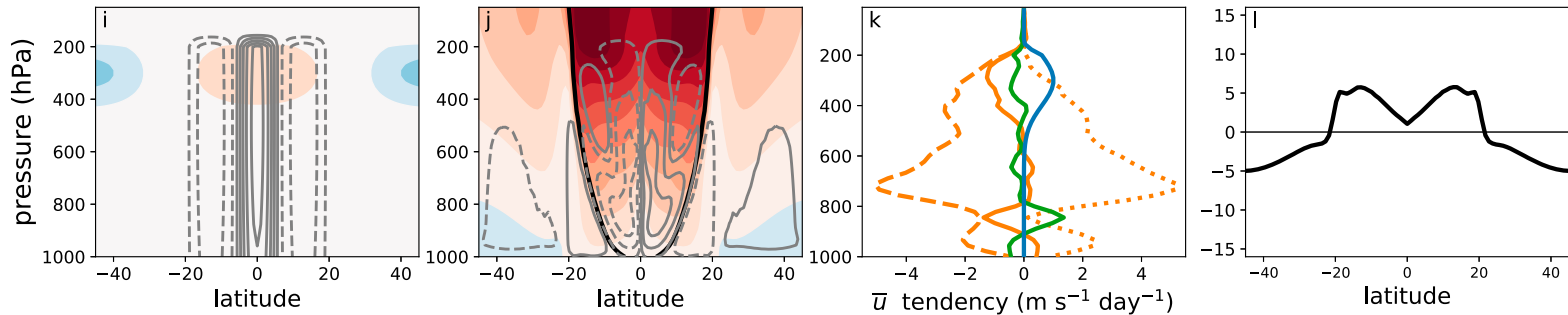
2



3



4



# Conclusions

- ▶ Elevated tropical temperatures lead to enhanced MJO-like variability
- ▶ MJO-like mode is trapped within the upper-level subtropical jets
- ▶ Momentum convergence is confined to upper troposphere and drives superrotation
- ▶ Sufficiently high vertical viscosity can bring superrotation down to the surface, but in GCM requires order-of-magnitude increase in CMT
- ▶ Transition to surface superrotation in past or future warm climate seems unlikely; on the other hand, understanding of CMT remains crude and surprises could be in store.



# Péclet number over broad range

$$Pe = \frac{\text{advective timescale}}{\text{diffusive timescale}} = \frac{\omega/p_s}{\nu/H^2}$$

