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## A story about convection, penetration, rotation and waves in stars

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*Warwick:* T. Goffrey

*Atlanta:* J. Pratt

- Brief introduction to physical conditions in stars (solar-like)
- Observational constraints on rotation, convection and waves in the Sun
- The first step toward a systematic numerical analysis of convection  $\longleftrightarrow$  penetration  $\longleftrightarrow$  waves  $\longleftrightarrow$  rotation  
*(work in progress)*

# Characteristics of stellar interiors and stellar convection

## - Complex regime of parameters

- $\text{Re} = \text{LU}/\nu \gg 10^6$  (*inertial/viscous*)  
turbulent convection
- Prandtl  $\text{Pr} \ll 1$

## - Characterised by very different timescales

**Sun**  $\tau_{\text{dyn}} \sim (R^3/GM)^{1/2} \sim 30 \text{ min}$   
 $\tau_{\text{conv}} \sim 6 \text{ days}$   
 $\tau_{\text{thermal}} \sim GM^2/(RL) \sim 2 \cdot 10^7 \text{ yr}$

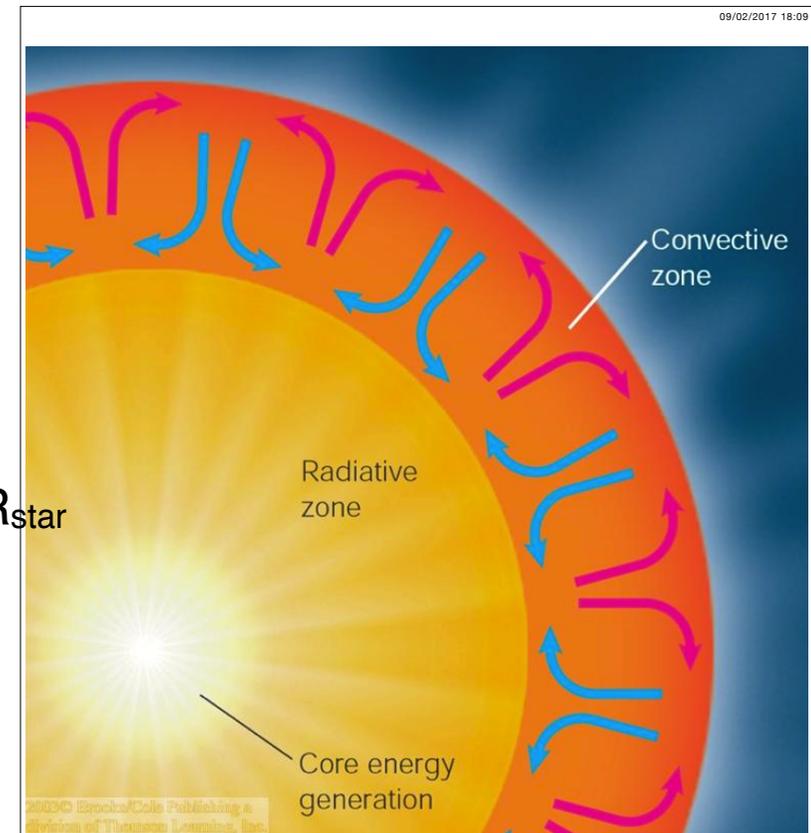
## - Very different lengthscales

Pressure scale height:  $H_P = dr/d\ln P$

centre:  $H_P \sim R_{\text{star}}$     Surface:  $H_P \sim 10^{-3} - 10^{-2} \times R_{\text{star}}$

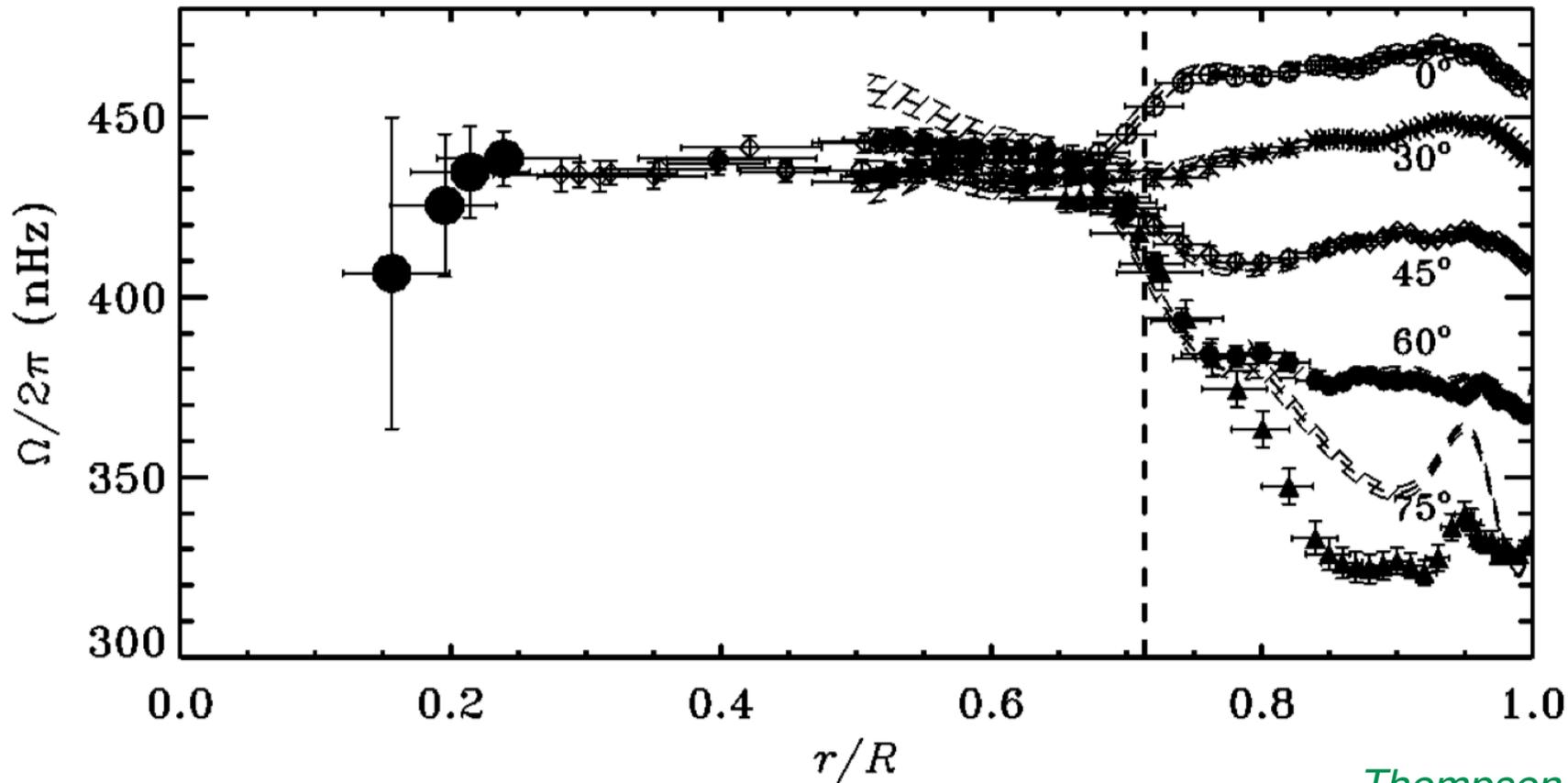
## - Range of Mach numbers ( $M \sim 10^{-6} - 1$ )

*surface: velocities*  $\sim c_{\text{sound}}$



# Solar rotation rate

From helioseismology: Inferred rotation rate as a function of fractional radius  $r/R$  and latitude



*Thompson et al. 2003*

- Solid body rotation in the stable radiative core
- Differential rotation in the convective envelope

☛ **Solar rotation: equator is faster than pole**

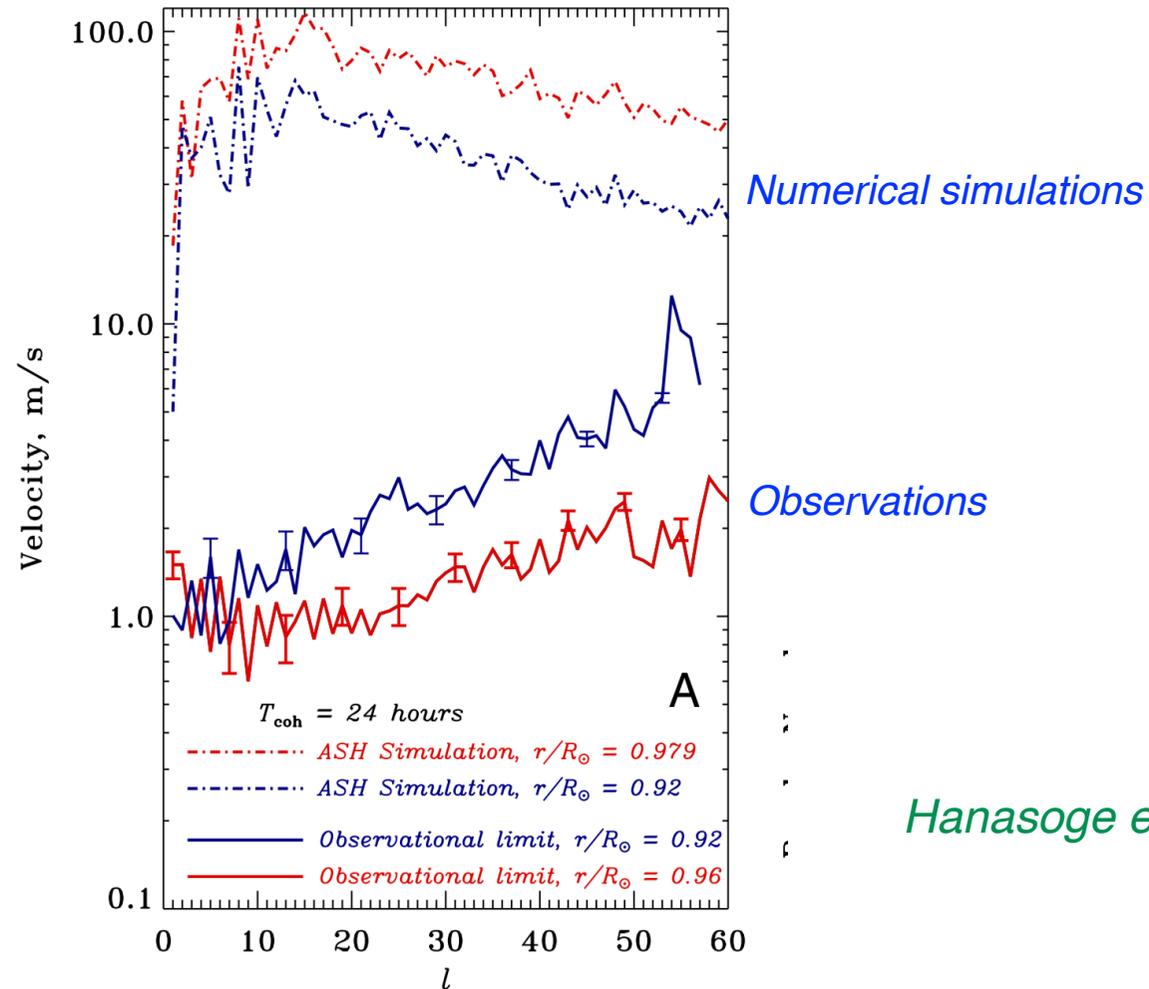
☛ The Sun's differential rotation is maintained by turbulent transport of angular momentum through the action of Reynolds stresses

☛ Spatial and temporal correlations between flow components that generate such stresses are thought to be imposed through Coriolis deflection of convective motions

⇒ **Very strong connection between convection and rotation**

# Solar subsurface convective velocities: the “convective conundrum”

Convective velocities as a function of spherical harmonic degree  $l$  and depth inferred from helioseismic data



*Hanasoge et al. 2012*

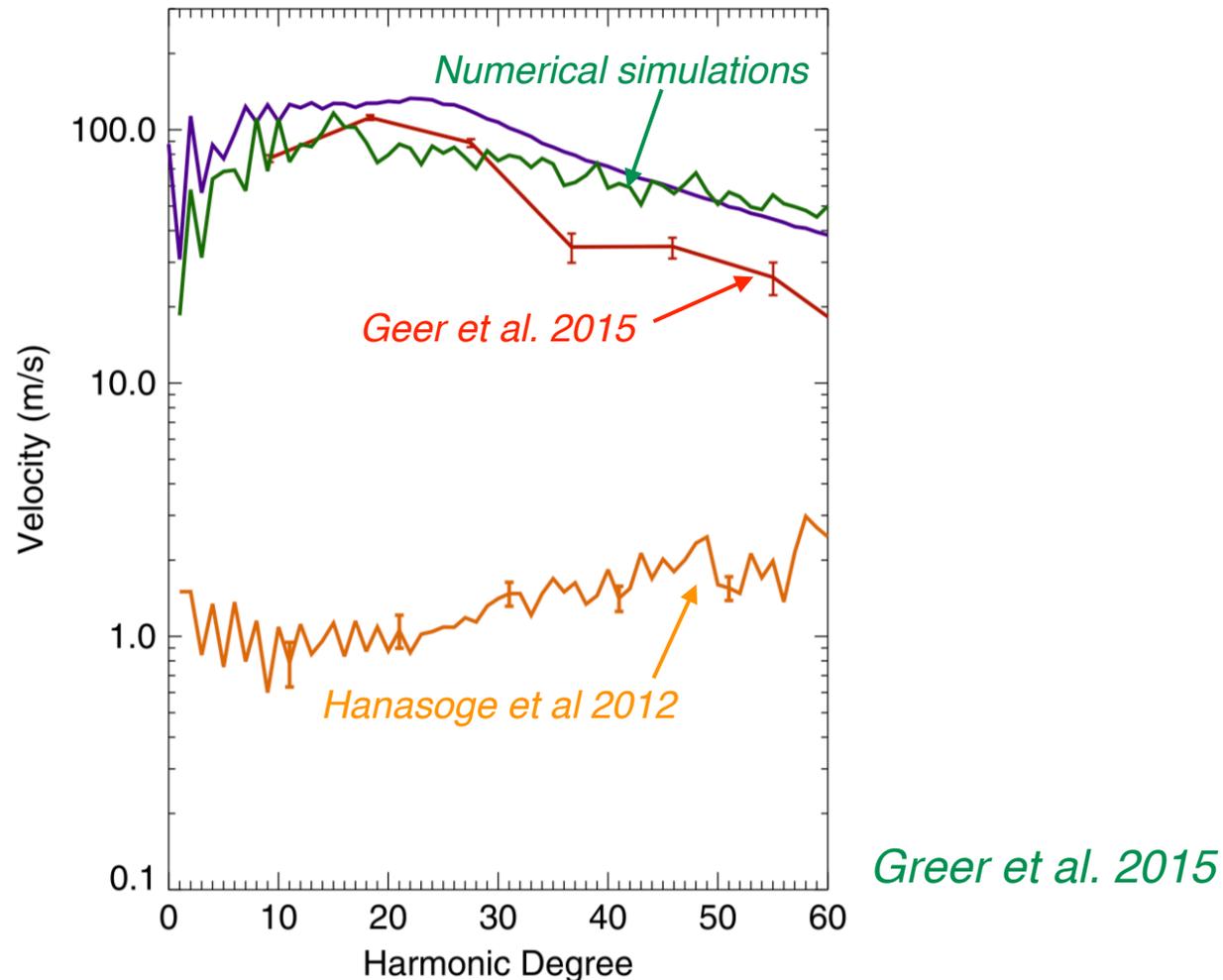
Large scale flows deep (depth  $\sim 30 \text{ Mm}$   $r/R_{\odot} \sim 0.957$ ) in the convective zone are significantly weaker than those predicted by most numerical simulations

→ Dynamical balance in numerical simulations is far from solar conditions?

Or

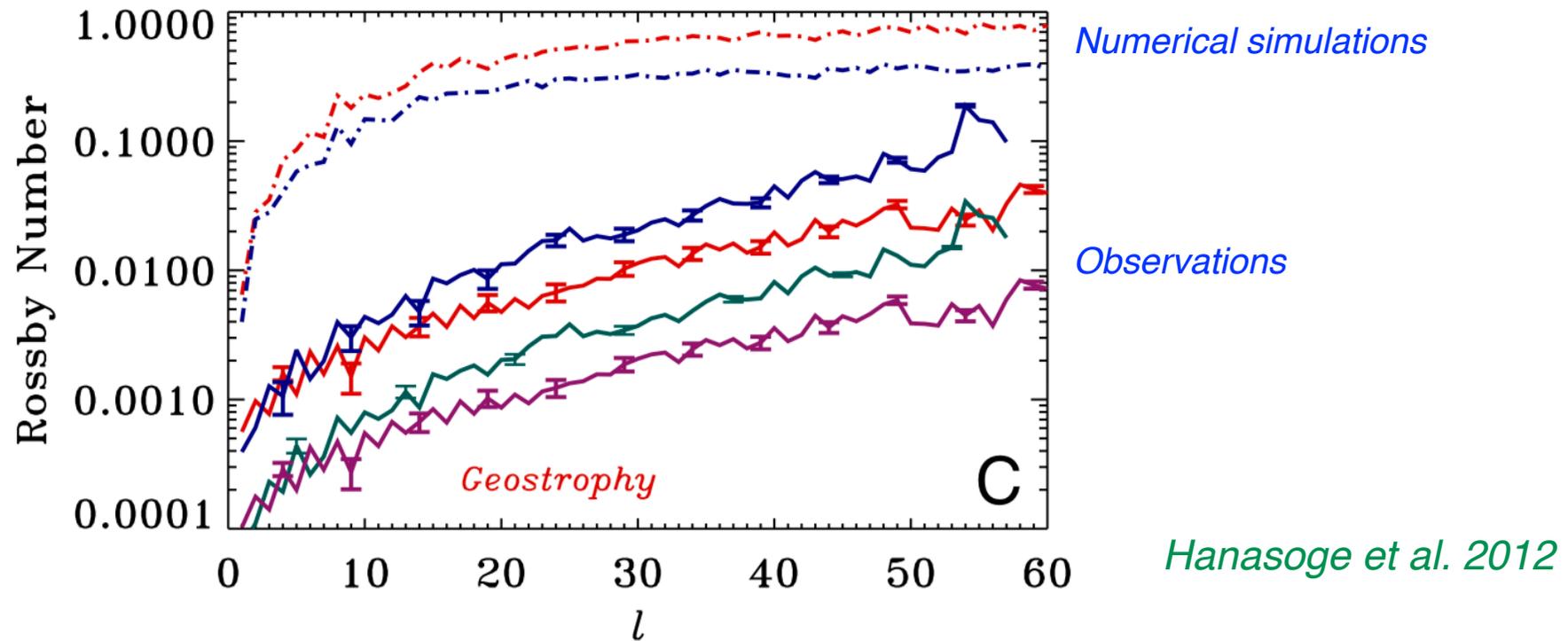
→ Problems with the method which underestimates velocities below few Mm?

Reanalysis of the solar data by Greer et al. 2015:



**The jury is still out.....**

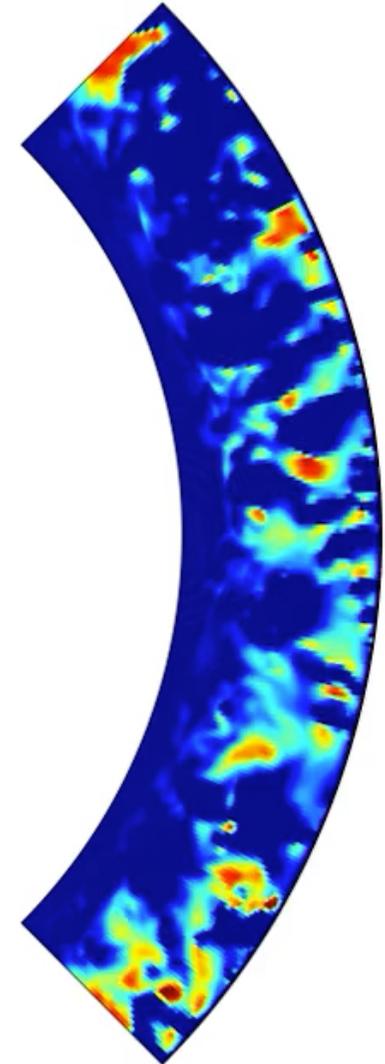
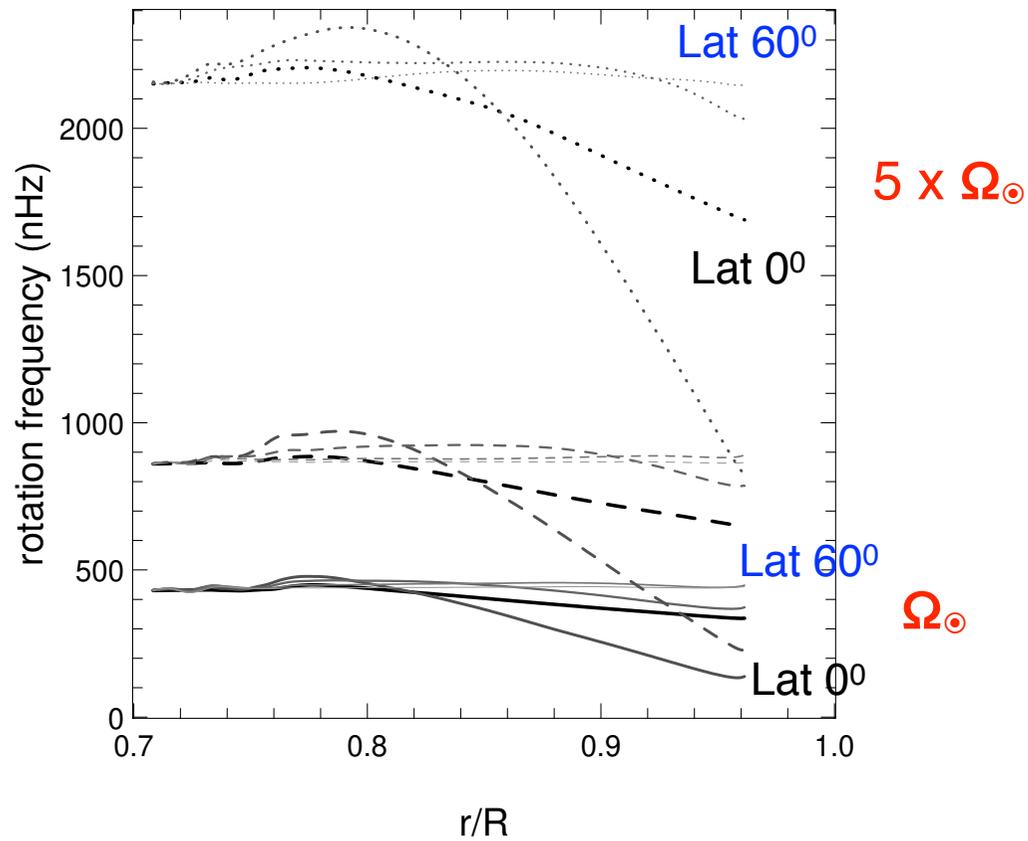
Observations suggest small Ro number at depth of  $\sim 30$  Mm and thus rotationally constrained convection



➡ According to Hanasoge et al. interior convection would be strongly geostrophically balanced (i.e., rotationally dominated) on these scales.

## Many numerical simulations are in the wrong rotational regime for solar-like conditions

3D simulations of solar convective envelope with our compressible code MUSIC  
( $r/R_{\odot} = 0.7 - 0.96$ )



*Constantino, Baraffe et al. in prep*

➤ Increased viscosity or lower diffusivity (**higher Pr**) or magnetic fields, can yield to solar-like rotation (and solve the solar convection conundrum)

➤ **Interesting:** convective penetration (overshooting) and inclusion of stable layers beneath the convective zone have a profound impact on differential rotation profile (*Beaudoin et al 2018, ApJ*)

➤ **Even more Interesting:** importance of taking into account non-local convection effects, i.e **surface cooling** leading to **downward plumes** that transport angular momentum radially inward  
→ despite high Pr, this leads to antisolar-rotation (*Karak et al 2018, PhFI*)

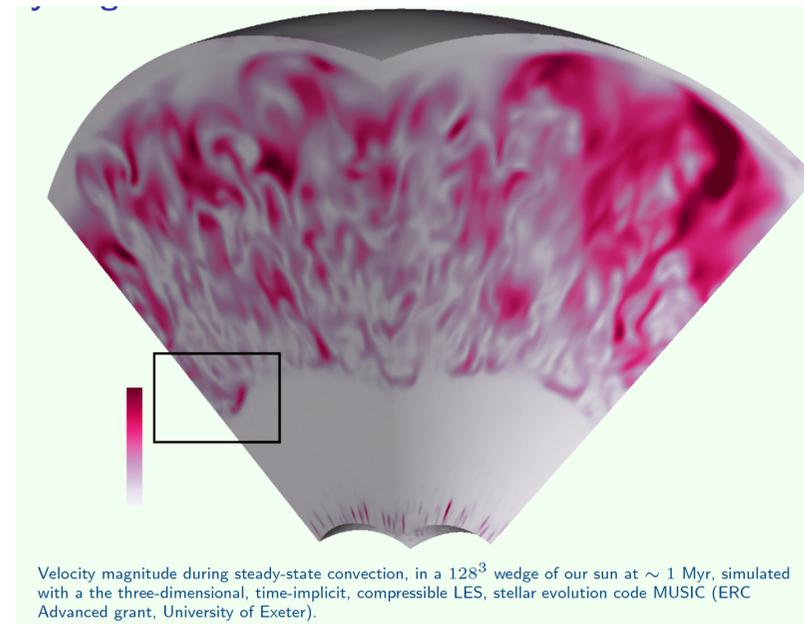
➤ **A consistent picture requires to account for surface cooling (near surface layers) and the interaction between convective zone/stable layers**

# Convective Penetration (“overshooting”)

**Penetration/Overshooting:** Extension of convective motions beyond the convective boundary in stably stratified regions

Boundary: Schwarzschild criterion for instability  $\frac{d \ln T}{d \ln P} > \left| \frac{d \ln T}{d \ln P} \right|_{\text{ad}}$

➔ Many observational evidences for presence of overshooting in stars of essentially all masses  $\geq 1 M_{\odot}$



*(Roxburgh 1965; Shaviv & Salpeter 1973; Schmitt et al 1984, etc...)*

## •Constraints from helioseismology

Frequencies of resonant acoustic waves depend on the adiabatic sound speed  $c$

$$c^2 = \Gamma_1 p/\rho \propto \Gamma_1 T/\mu$$

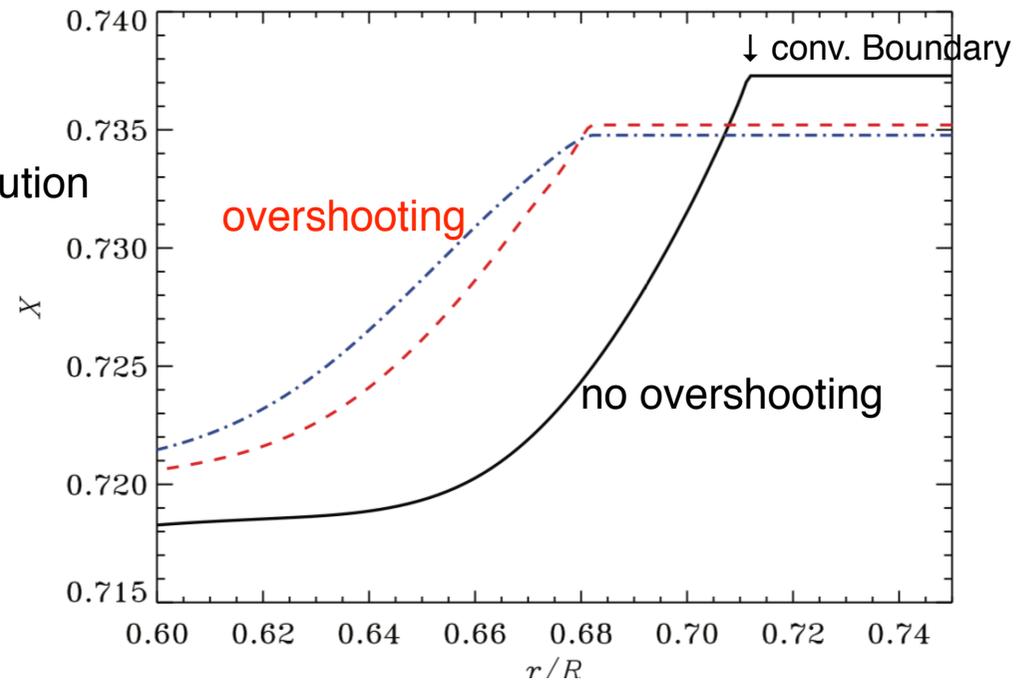
☛ helioseismology provides information about the profile of  $c^2$  in the Sun's interior

⇒ Indication of **extra mixing below the convective boundary**

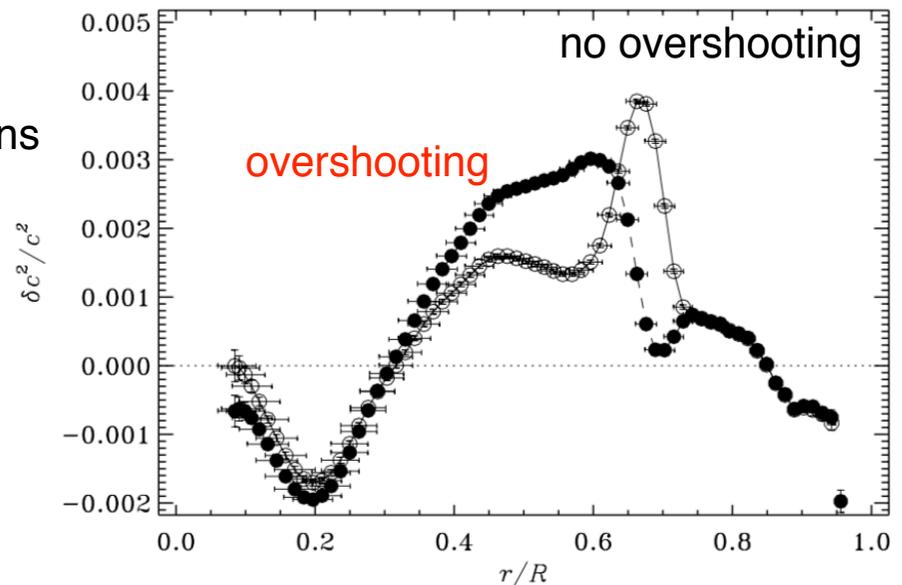
*Christensen-Dalsgaard 1993*

➡ Convective penetration most likely the best process driving chemical mixing

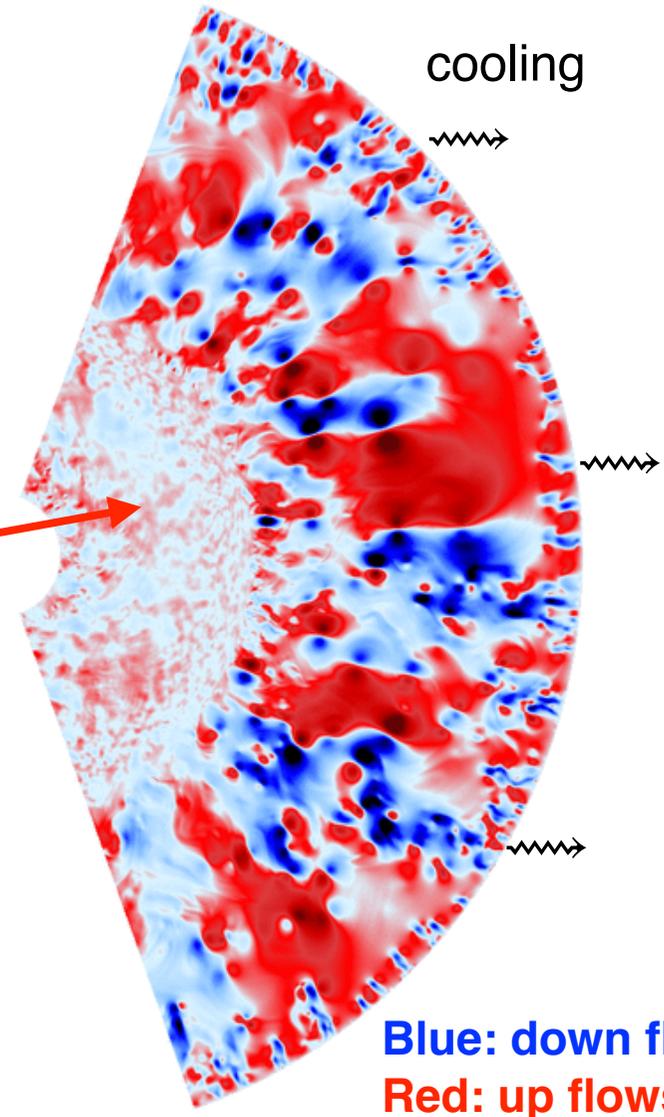
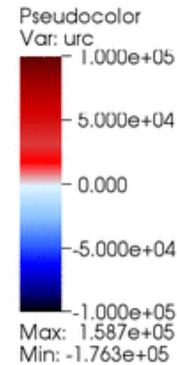
Profile of hydrogen abundance close to the lower convective boundary from 1D stellar evolution models (phenomenological approach)



➡ Relative difference  $\delta c^2/c^2$  between Sun observations and model



# Internal Waves (gravity or gravito-inertial)



## Internal waves excited by

- **Reynolds stresses** or **entropy fluctuation** due to convection near the bottom of the convective zone
- and/or **penetrative plumes** just below the boundary

*Press 1981; Schatzman 1993; Zahn et al. 1997;  
Kumar & Quataert 1998; Goldreich et al. 1994, etc...*

**Blue: down flows**  
**Red: up flows**

*Radial Velocity magnitude: characteristic patterns of convection with up-wards and down-wards "plumes"*

↳ **Transport of angular momentum** by convectively (and/or penetrating flows) excited internal waves

- Suggestion of a similar mechanism as the QBO in the Sun with shear-layer oscillations in the “tachocline” (Kumar et al. 1999)

Major mechanism for the wave damping in stellar interior is radiative damping:

Damping distance:  $d \propto \frac{w^4}{l^3 K_T}$       Thermal diffusivity  $K_T = \frac{16\sigma T^3}{3\rho^2 \kappa c_p}$

⇒ **larger damping (i.e small d) for decreasing frequency  $w$**

## Basic picture:

Frequency of waves excited at convective boundary  $r_c$  in the rest frame at radius  $r$ :

$$\omega_{\text{rf}}^m(r) = \omega_c + m [\Omega_c - \Omega(r)]$$

If  $\Omega(r)$  increases with depth below the convective zone ( $\Omega(r) > \Omega_c$ )

→ **Prograde** ( $m > 0$ ) are shifted to **lower**  $\omega$

→ **Retrograde** ( $m < 0$ ) are shifted to **higher**  $\omega$

⇒ **Damping of prograde waves is enhanced (shorter damping length) relative to retrograde waves ( $d \propto \omega^4$ )**

⇒ positive angular momentum deposited just below  $r_c$  while retrograde waves deposit their negative angular momentum further away

➡ **double-peaked shear layer**

⇒ prograde shear layer propagates toward  $r_c$  to eventually merge with shear layer just below convective zone

☞ process repeats leading to oscillatory behaviour i.e QBO-like process

Attempts to find evidence of this process in hydrodynamical simulations of the Sun

- No QBO-like oscillations at the convective boundary (tachocline)  
Suggestion of a low-amplitude QBO-like oscillation well below the convective boundary  
*(Rogers et al. 2006)*
- Simplified numerical experiments can produce oscillating shear flows  
*(Rogers et al. 2008)*
  - ☞ Importance of critical layers where waves are attenuated and which could dominate the mean flow dynamics
  - ☞ **Importance of the spectrum and amplitudes** of waves driven (i.e excitation mechanism of the waves → convection/turbulence/plumes?)

Press 1981: “*Convective overshoot and sub scale turbulent instability of a highly nonlinear internal wave are two aspects of a single, but **messy**, physical process*”

To understand/describe waves and transport of angular momentum in stellar radiative cores  
↪ **one needs to understand the full connection(s) between convection-penetration-rotation-waves**

**A very long journey!**

Our first step of this thousand miles journey is to **understand convection and penetration** and their close inter-connection (*which impact convective velocities, differential rotation, waves and vice-versa....*)

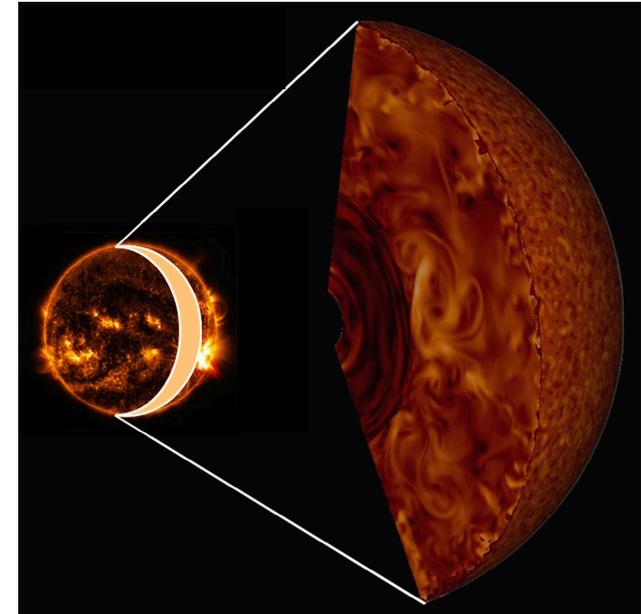
## Our approach of the problem: a systematic survey of convection/penetration based on 2D and 3D fully compressible time implicit simulations

### *In progress*

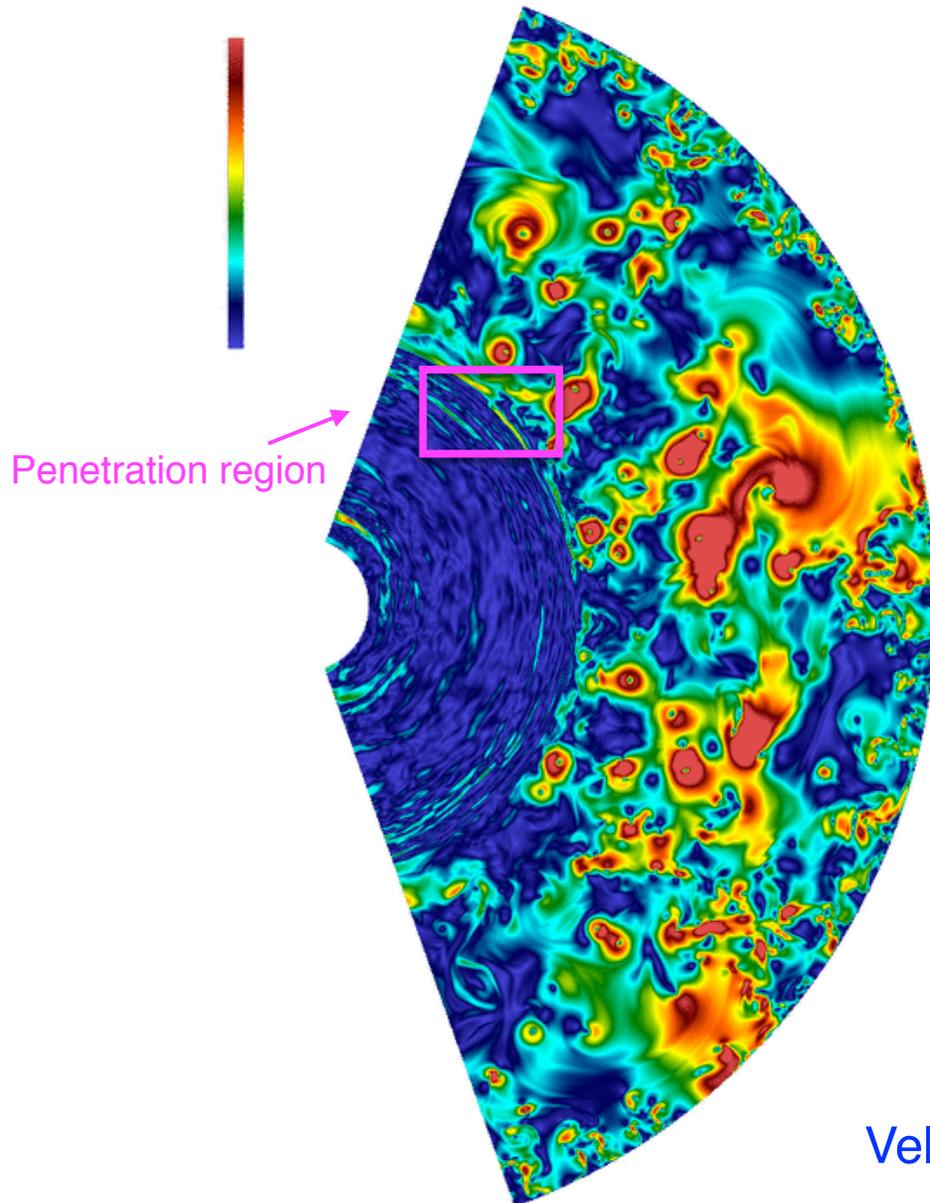
- Explore the impact of radial extension of the numerical domain:  
Near-surface layers + stable radiative layers
- Analyse the excitation of waves

### *Next steps:*

- Explore the impact of rotation and transport of angular momentum
- Explore the impact of magnetic fields



- **Primary goal:** understand the connection between the dynamics in the convective zone and plumes in the penetration region (characterise the mixing, wave excitation, etc..)



Numerical simulations with MUSIC

(*Viallet et al. 2011, 2013, 2016; Goffrey et al. 2017*)

- Spherical geometry (2D or 3D)
- Finite volume scheme
- **Time implicit solver**
- **Fully compressible** hydrodynamics
- Thermal diffusion  $\nabla \cdot (\chi \nabla T)$

*Radiative conductivity*  $\chi = 16\sigma T^3 / 3\kappa\rho$

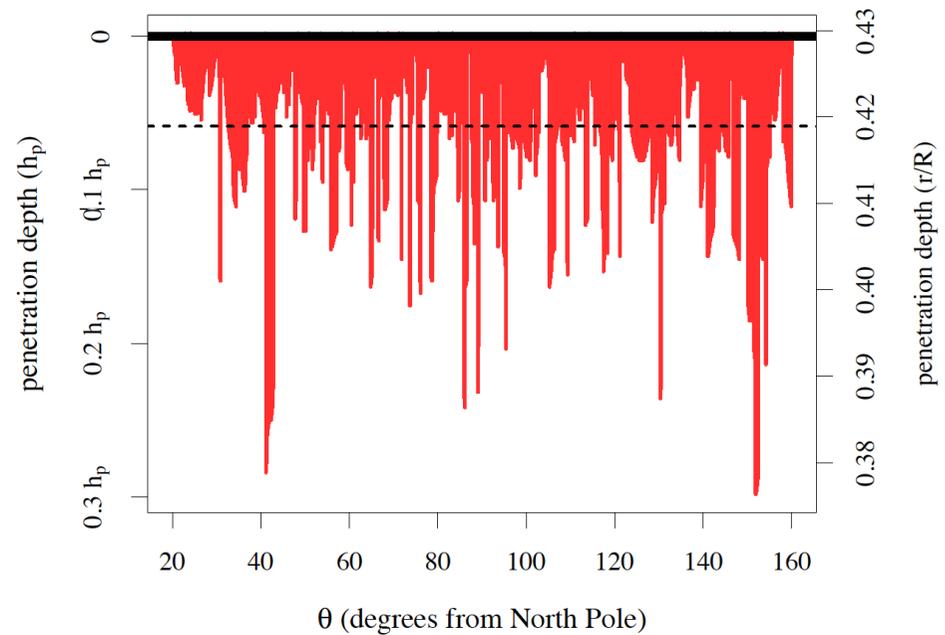
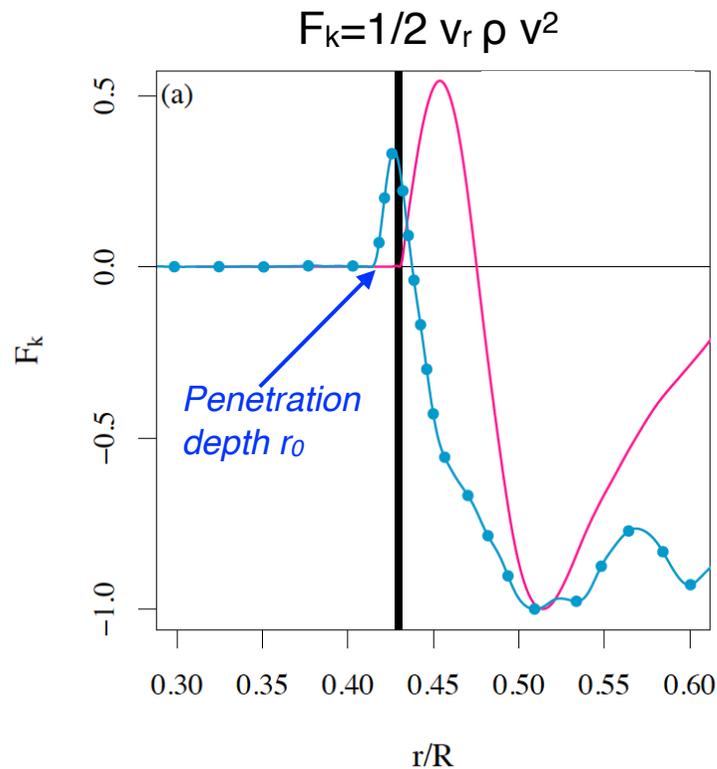
- Realistic stellar stratification

*Realistic equation of state (ionisation, partial degeneracy, mixture of composition, etc)*

Velocity magnitude : very high res 2432x2048

## ➤ A new approach to describe convective penetration: Extreme value statistics (Pratt et al. 2017, 2019)

Typical shape of the penetration depths (at a given time): extent of downflows (defined by vertical flux of kinetic energy  $\rightarrow 0$ ) beyond the convective boundary varies with colatitude  $\theta$

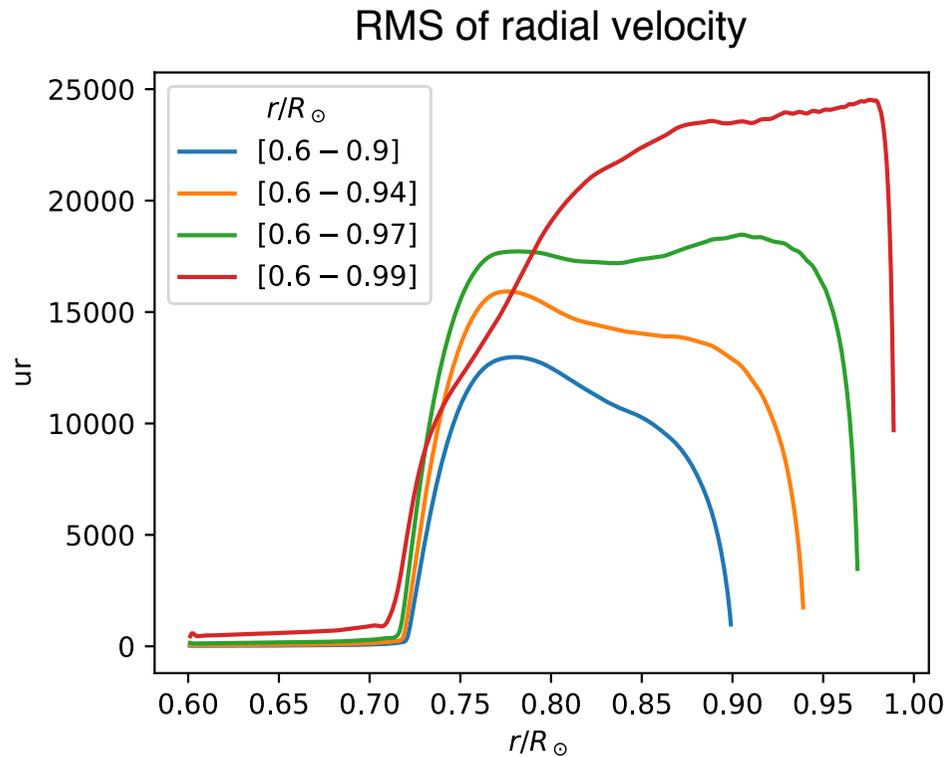


PDF of penetration depths  $r_0$  for all downflows (over the whole simulation time) and the PDF of maximal penetration depths  $r_{max}$

## Impact of radial extension of the numerical domain:

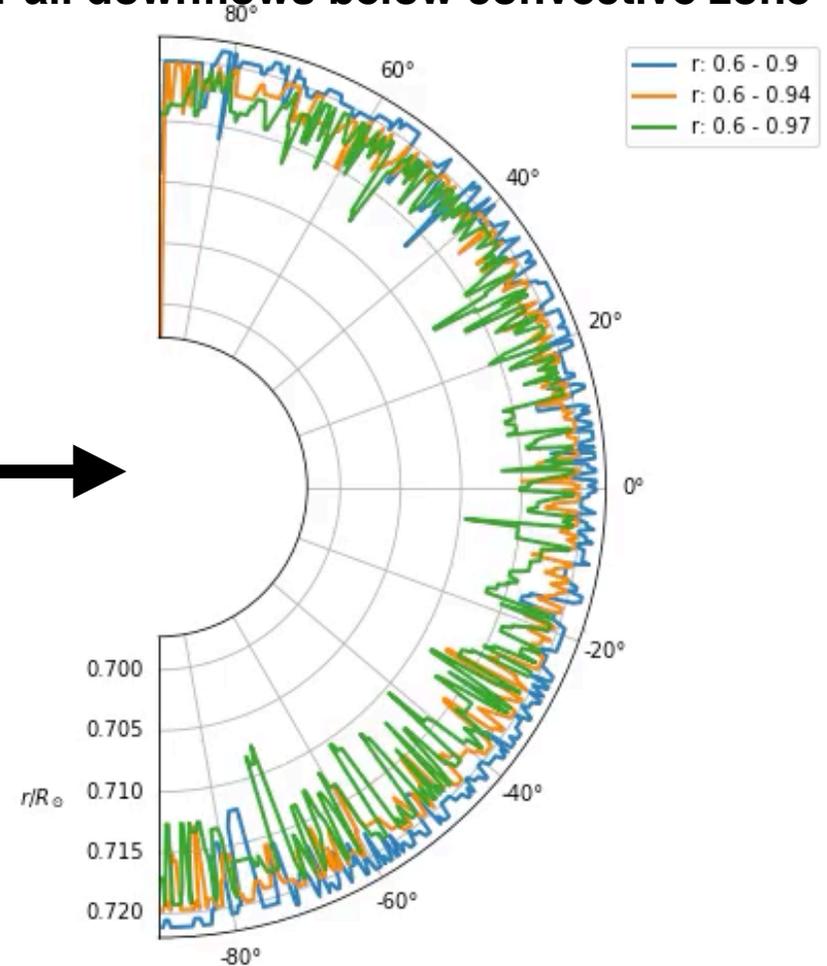
(*In progress*) wide survey (2D/3D) of a solar-like model with variable radial extensions in the range  $r/R = 0.1 - 1$

- General **increase** of radial velocities when including more near-surface layers (effect of **surface cooling** on the **strength of convective plumes**)

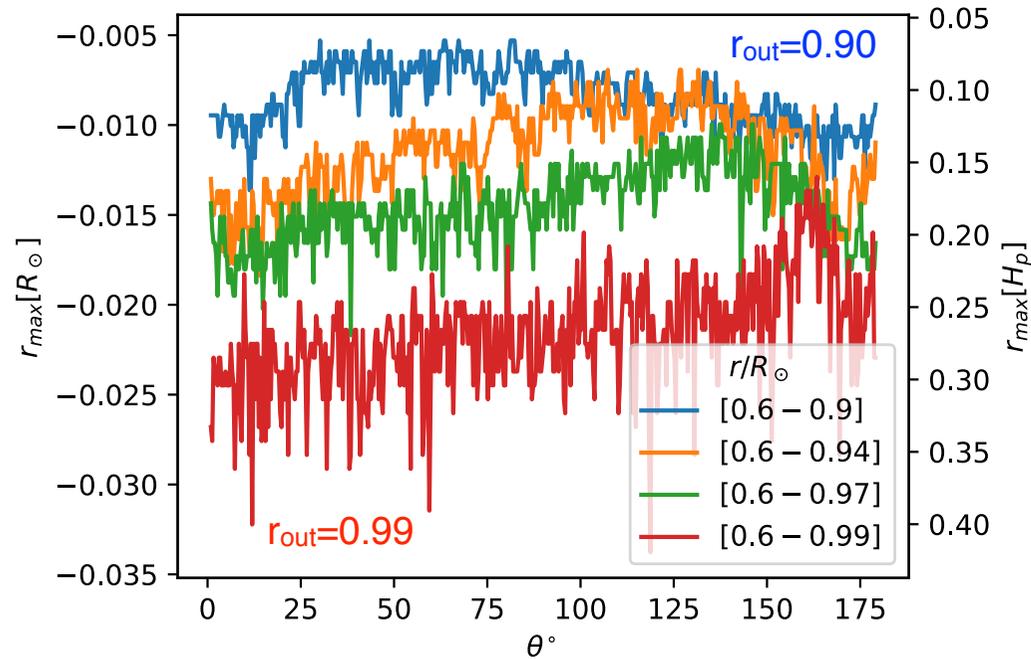


*Vlaykov et al., in prep*

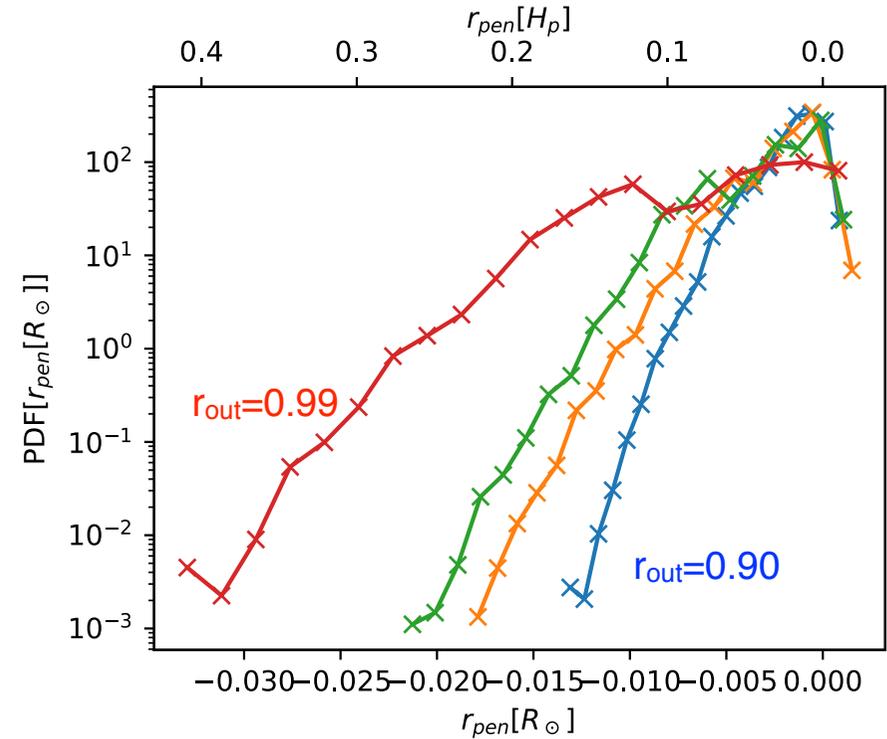
## Determination of the Penetration depths $r_0$ for all downflows below convective zone



## Maximum penetration depths $r_{\max}$ below convective zone



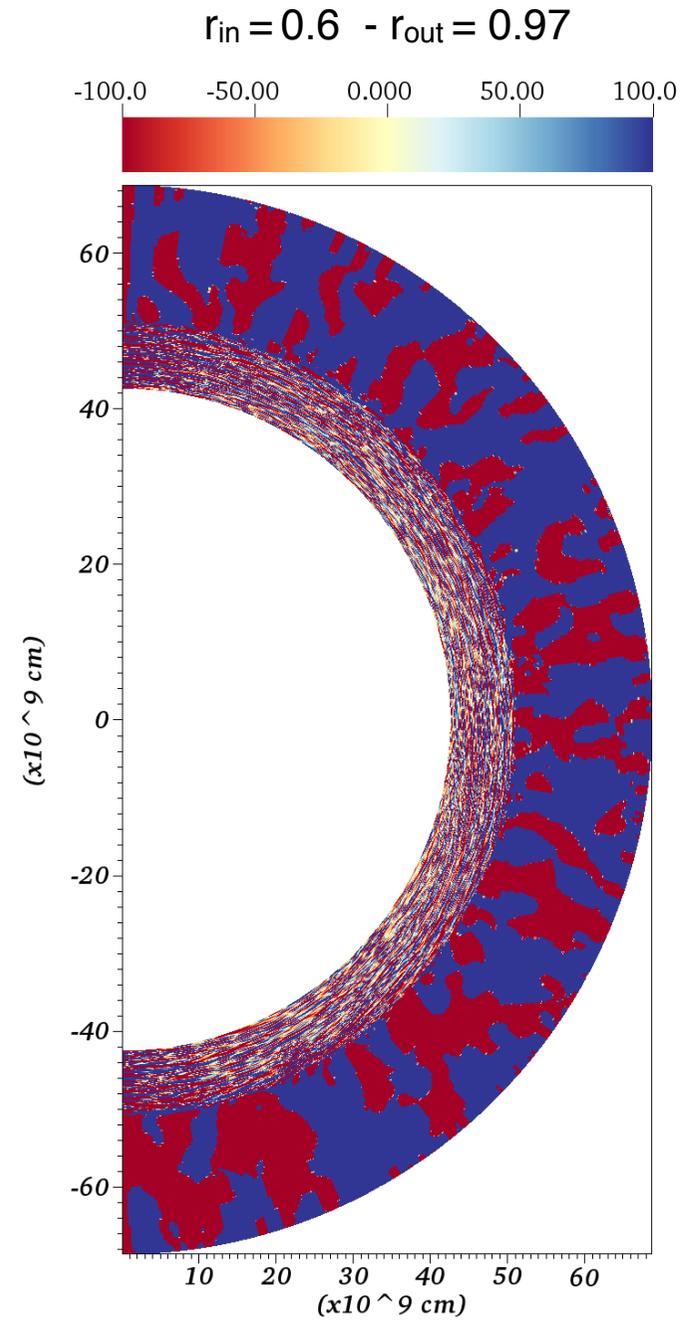
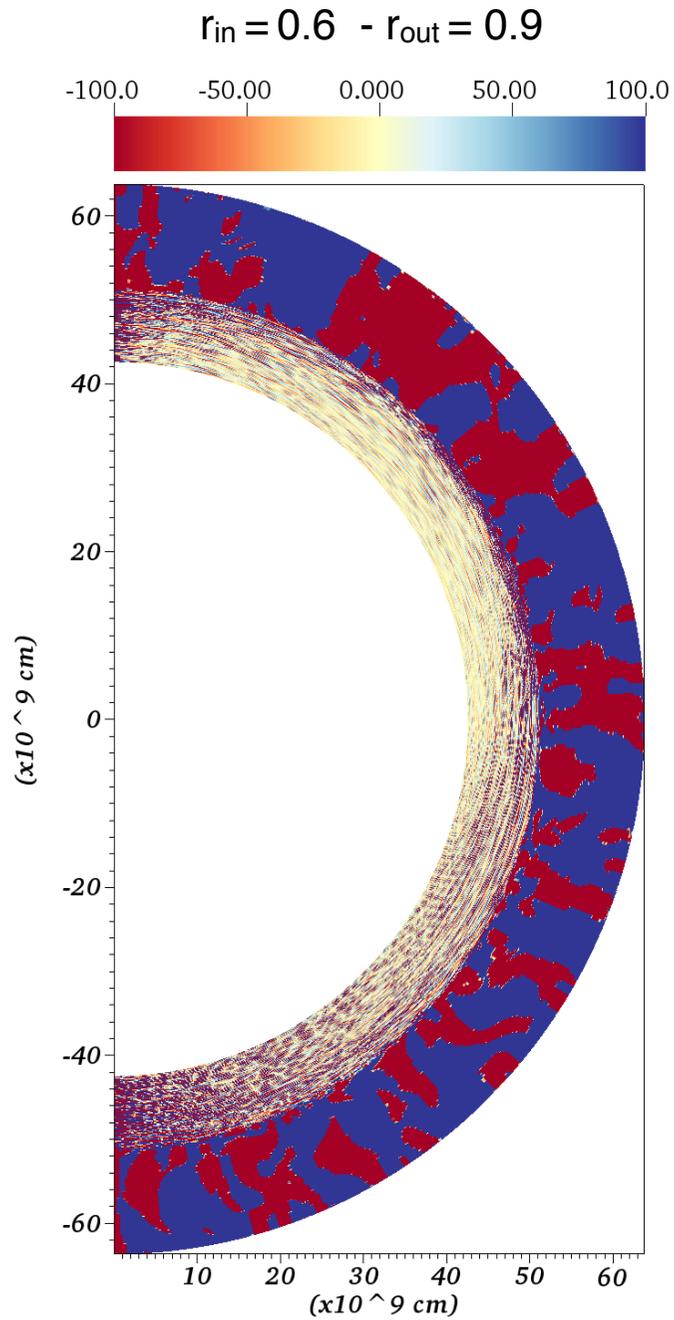
## PDF of penetration depths $r_0$



➡ The **closer** the upper domain to the **near-surface layers** where **cooling** takes place, the **stronger the downwards plumes** and the **deeper the penetration plumes**

➡ Ongoing analysis of the impact of including a larger inner domain (stably stratified region extending down to  $r_{\text{in}}=0.1$ )

## Radial velocity snapshot for two different outer extensions



➤ **Next step: Quantitative analysis of internal waves**

Radial velocity fields are used to characterise internal gravity waves that propagate in the stable zone

From the radial velocity fields:  $v_r(r, \theta, \varphi, t)$

—> projection on spherical harmonic basis  $Y_l^m$ .  $\mathbf{e}_r$

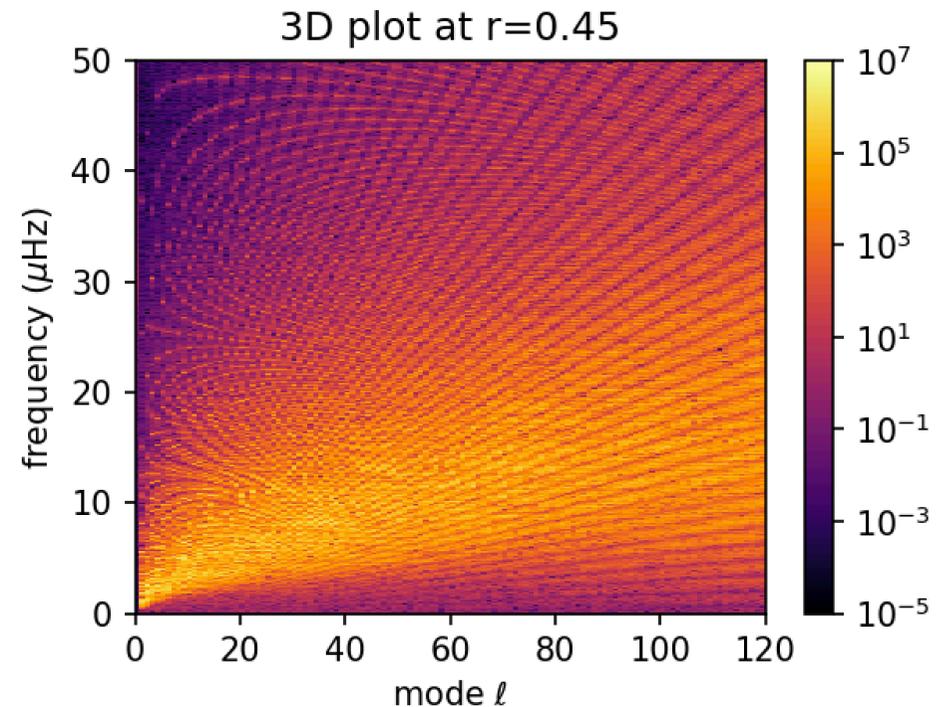
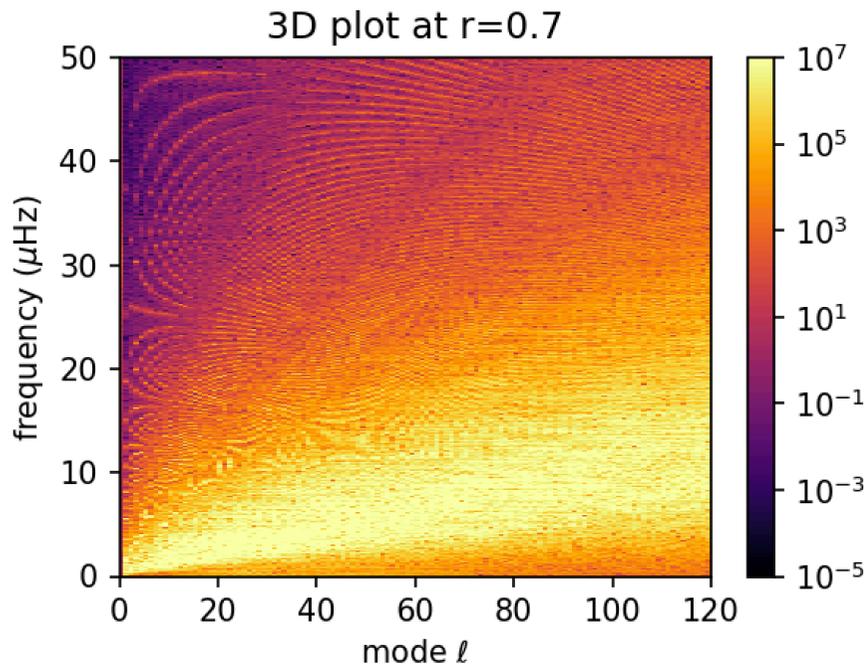
—>  $v_r'(r, l, m, t)$

—> Fourier transform

—>  $v_r''(r, l, m, \omega)$  ( $\omega$  is the mode frequency)

Illustration for a Solar model  $r_{in}=0.4$  -  $r_{out}=0.9$

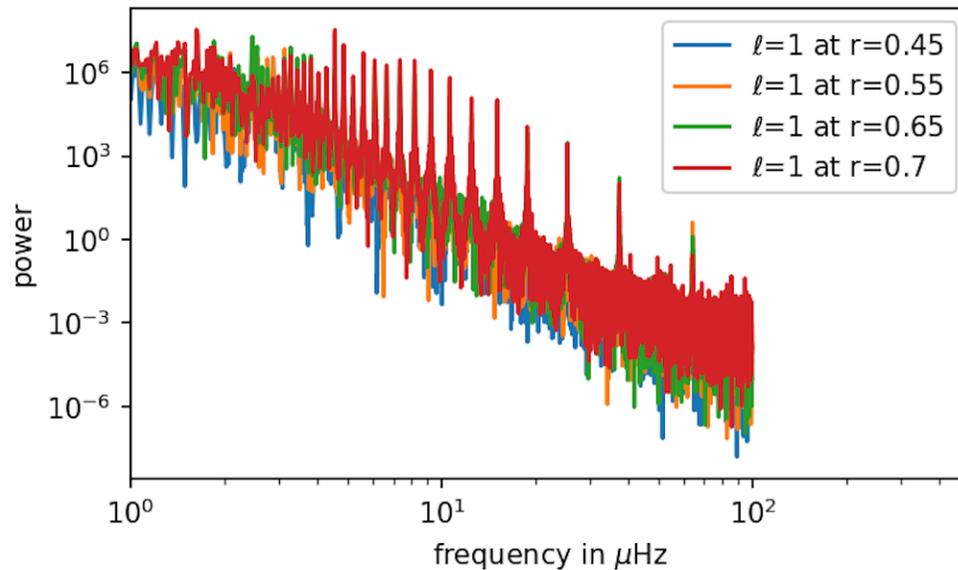
Spectrum of waves at different depths:  $(v_r'')^2$



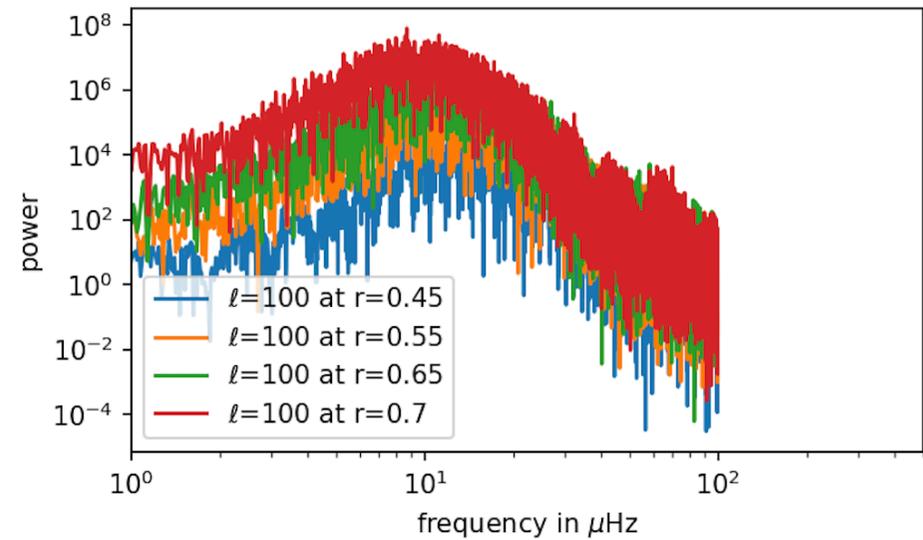
*Lesaux et al.*

# Spectrum for different $l$ at different depths

$l=1$



$l=100$



- ➡ Powerful tool to analyse the excitation spectrum of waves (stresses due to convective eddies at the boundary and/or **penetrating plumes**)

*Rogers & MacGregor 2011; Alvan et al. 2015*

- ➡ A motivation: Impact of extreme penetration events?

**Story to be continued....**

## Conclusion (or thoughts...)

- Solar rotation, convective velocities, waves and convective penetration are linked  
⇒ **Need for a global framework**
- **Surface cooling and plumes dynamics play an important role**
- **What is the role of extreme penetrating plume events on wave excitation mechanism and properties?**